HWY SOLUTIONS
(1) $K_{B E A M}=K_{B}=\frac{A E}{L}$
$k_{\text {lect }}$ for $K_{B} \& K_{1}$

$$
k_{\text {est }} \frac{1}{\left(\frac{1}{k_{B}}+\frac{1}{k_{1}}\right)}=\frac{k_{1} k_{B}}{k_{1}+k_{B}}
$$


$k_{\text {eft }}$ \& $k_{2}$ act in parallel

$$
\begin{aligned}
k_{\text {eff }}=k_{\text {left }}+k_{2} & =k_{2}+\frac{k_{1} A E}{L\left(k_{1}+\frac{A E}{L}\right)} \\
k_{\text {eff }} & =k_{2}+\frac{k_{1} A E}{L k_{1}+A E}
\end{aligned}
$$



Sos Eqn

$$
\begin{aligned}
& m \ddot{x}+k_{\text {eff }} x=0 \\
& \ddot{x}+\frac{k_{\text {eff }} x}{m}=0 \\
& w_{\Lambda}=\sqrt{\frac{k_{\text {eff }}}{m}}
\end{aligned}
$$

(2) (a) Draw system as


$$
\begin{gathered}
f_{k 1}=k a \theta \\
\left(m_{\text {bean }}\right)_{p}=\left(I_{\text {beam }}\right)_{p} \ddot{\theta} \\
f_{k 2}=-k \cdot 2 a \cdot \theta \\
f_{M}=M 3 a \ddot{\theta}
\end{gathered}
$$

Positive $\theta$ CW 5

$$
\begin{aligned}
& \sum M_{p}=0=-f_{k 1} \cdot a-\left(M_{\text {beam }}\right)_{p}^{+} f_{k 2} \cdot 2 a-f_{m} \cdot 3 a \\
& 0=-k a^{k} \theta-\frac{7}{3} m a^{2} \ddot{\theta}-k \cdot 4 a^{2} \cdot \theta-9 a^{2} M \cdot \ddot{\theta} \\
& \left(\frac{7}{3} m+9 m\right) \ddot{\theta}+(5 k) \theta=0 \quad\left(I_{\text {beam }}\right)_{p}=\underbrace{\left(I_{\text {beam }}\right)_{c m}+m a^{2}}_{\text {parallel axis }} \text { theorem } \\
& \ddot{\theta}+\frac{5 k}{\frac{5 k}{3} m+9 m} \theta=0 \quad\left(I_{\text {beam }}\right)_{c m}=\frac{1}{12} m \cdot(4 a)^{2}
\end{aligned}
$$

| (b) $T$ | $V$ |  |
| :---: | :---: | :---: |
| beam | $\frac{1}{2}\left(I_{b}\right)_{p} \dot{\theta}^{2}$ | 0 |
| left $k$ | 0 | $\frac{1}{2} k(a \theta)^{2}$ |
| right $k$ | 0 | $\frac{1}{2} k(2 a \theta)^{2}$ |
| mass $M$ | $\frac{1}{2} M(3 a \dot{\theta})^{2}$ | 0 |

$$
\left(I_{b}\right)_{p}=\left(I_{\text {beam }}\right)_{p}=\frac{7}{3} m a^{2}
$$

$$
\begin{aligned}
L=T-V \quad L & =\frac{1}{2}\left(I_{b}\right)_{p} \dot{\theta}^{2}+\frac{9}{2} a^{2} m \dot{\theta}^{2}-\frac{1}{2} k a^{2} \theta^{2}-\frac{1}{2} k 4 a^{2} \theta^{2} \\
L & =\left(\frac{1}{2}\left(I_{b}\right)_{p}+\frac{9}{2} a^{2} M\right) \dot{\theta}^{2}-\left(\frac{5}{2} k a^{2}\right) \theta^{2}
\end{aligned}
$$

$$
\begin{aligned}
& \frac{d}{d t} \frac{\partial L}{\partial \dot{\theta}}-\frac{\partial L}{\partial \theta}=0 \\
& \frac{d}{d t}\left[2\left(\frac{1}{2}(I b)_{p}+\frac{9}{2} a^{2} m\right) \dot{\theta}\right]+3 k a^{2} \theta=0 \\
& \left((I b)_{p}+9 a^{2} m\right) \ddot{\theta}+5 k a^{2} \theta=0 \\
& \left(\frac{7}{3} m q^{2 x}+9 m q^{2}\right) \ddot{\theta}+5 k q^{2} \theta=0 \\
& \ddot{\theta}+\frac{5 k}{\left(\frac{7}{3} m+9 m\right)} \theta=0
\end{aligned}
$$

Note: For KE of beam, can alter natvel, use

$$
\begin{aligned}
& T_{\text {beam }}=\frac{1}{2}\left(I_{\text {beam }}\right)_{c m} \dot{\theta}^{2}+\frac{1}{2} m(a \dot{\theta})^{2} \\
& \text { as K.E. rigid body }=\frac{1}{2} I_{c m} \dot{\theta}^{2}+\frac{1}{2} m V_{c m}^{2}
\end{aligned}
$$

Keep in mind both rotational \& linear motion
(3) Use small \& approx. Do not neglect gravity.


PE loss by rotation of beam by $\& \theta$
mass M: Mig. $\Delta h_{M}$ beam! $m \cdot g \cdot \Delta h_{m}$

$$
\begin{aligned}
& \Delta h_{M}=(3 \cdot a) \cdot(1-\cos \theta) \\
& \Delta h_{m}=\left(\frac{3}{2} a\right) \cdot(1-\cos \theta) \\
& s m_{a l l} \& \rightarrow \cos \theta \simeq 1-\frac{\theta^{2}}{2} \\
& \Delta h_{M}=3 a \frac{\theta^{2}}{2} \quad \Delta h_{m}=\frac{3}{2} a \frac{\theta^{2}}{2}
\end{aligned}
$$

$$
\begin{aligned}
& L=T-V \\
& \begin{aligned}
& L=\frac{9}{2} m a^{2} \dot{\theta}^{2}+\frac{q}{24} m a^{2} \dot{\theta}^{2}+\frac{1}{2} m\left(\frac{9}{4}\right) a^{2} \dot{\theta}^{2}+M g 3 a \frac{\theta^{2}}{2}+m g \frac{3}{2} a \frac{\theta^{2}}{2} \\
&-4 K a^{2} \theta^{2}
\end{aligned}
\end{aligned}
$$

Lagrange

$$
\begin{aligned}
& \quad \frac{d}{\partial t} \frac{\partial L}{\partial \dot{\theta}}-\frac{\partial L}{\partial \theta}=0 \\
& 0=\frac{d}{d t}\left[\left(9 m a^{\alpha}+\frac{3}{4} m a^{\not x}+\frac{9}{4} m a^{\alpha x}\right) \dot{\theta}\right]-\left(m g 3 \alpha \alpha+m g^{\frac{3}{2}} \alpha \alpha-8 k a^{\alpha}\right) \theta \\
& 0=(9 m a+3 m a) \ddot{\theta}+\left(-m g 3-m g \frac{3}{2}+8 k a\right) \theta
\end{aligned}
$$

Effective Stiffness is coefficient in front of $\theta$ (as in $m \ddot{x}+k x=0$ )
system is unstable (by divergence) if effective stiffness $<0$
Stability:

$$
\begin{aligned}
& 8 k a-\left(3 m g+\frac{3}{2} m g\right)>0 \\
& 8 k a>3 m g+\frac{3}{2} m g \\
& k>\frac{\left(3 m g+\frac{3}{2} m g\right)}{8 a}
\end{aligned}
$$

In order to see how the $x=c_{1} \cos \omega t+c_{2} \sin \omega t$ solution agrees with the $x=1+t$ solution in the limit as $\omega$ approaches zero, we observe that the initial conditions associated with $x=1+t$ are:

$$
\left\{\begin{array}{l}
x(0)=1 \\
x^{\prime}(0)=1
\end{array}\right.
$$

applying these initial conditions to $x=c_{1} \cos \omega t+c_{2} \sin \omega t$, we get:

$$
\left\{\begin{array} { l } 
{ x ( 0 ) = 1 } \\
{ x ^ { \prime } ( 0 ) = 1 }
\end{array} \Rightarrow \left\{\begin{array}{l}
c_{1}=1 \\
\omega c_{2}=1 \Rightarrow c_{2}=\frac{1}{\omega}
\end{array}\right.\right.
$$

so the solution becomes:

$$
x=\cos \omega t+\frac{1}{\omega} \sin \omega t
$$

now, we take the limit of this solution as $\omega \rightarrow 0$ :

$$
\lim _{\omega \rightarrow 0} x=\lim _{\omega \rightarrow 0} \cos \omega t+\lim _{\omega \rightarrow 0} \frac{1}{\omega} \sin \omega t=1+\lim _{\omega \rightarrow 0} \frac{t \cos \omega t}{1}=1+t
$$

where we use l'hopital's rule to evaluate the second limit.
In summary, since one of the arbitrary constants actually depends on $\omega$, it is essential to account for it when taking the limit $\omega \rightarrow 0$.

在 1.



$$
\begin{aligned}
& y^{2}+(d-x)^{2}=a^{2} \\
& y=\sqrt{a^{2}-(d-x)^{2}}
\end{aligned}
$$

$$
\begin{aligned}
& T=\frac{1}{2} m_{1} \dot{x}^{2}+\frac{1}{2} m_{2} \dot{y}^{2}, V \\
& 2 y \dot{y}-2(d-x) \dot{x}=0 \Rightarrow \dot{y}=\frac{d-x}{2} x^{2} \\
& \sqrt{a^{2}-(d-x)^{2}} \dot{x} \\
& \therefore T=\frac{1}{2} f(x) \dot{x}^{2}, f(x)=m_{1}+\frac{(d-x)^{2} m_{2}}{a^{2}-(d-x)^{2}} \\
& \mathcal{L}=T-V=\frac{1}{2} f(x) \dot{x}^{2}-\frac{1}{2} k x^{2} \\
& \frac{\partial \mathcal{L}}{\partial x}=\frac{1}{2} f^{\prime} \dot{x}^{2}-k x \\
& \frac{\partial \mathcal{L}}{\partial \dot{x}}=f \dot{x}, \frac{d}{d t} \frac{\partial f}{\partial \dot{x}}=f \ddot{x}+f^{\prime} \dot{x}^{2} \\
& \frac{d}{d t} \frac{\partial \mathcal{L}}{\partial \dot{x}}-\frac{\partial f}{\partial x}=0 \Rightarrow f \ddot{x}+f^{\prime} \dot{x}^{2}-\frac{f^{\prime} \dot{x}^{2}}{2}+k x
\end{aligned}
$$

\#1 continued
So Lagrange's equation is

$$
f \ddot{x}^{\prime}+f^{\prime} \frac{x^{2}}{2}+k x=0 \quad \text { Ansitopart a }
$$

where $f(x)=m_{1}+\frac{(d-x)^{2} m_{2}}{a^{2}-(d-x)^{2}}$
b) To find national frequency, beep only lineiac tums. This gives

$$
\begin{gathered}
f(0) \ddot{x}+k x=0 \\
w^{2}=\frac{k}{f(0)}=\frac{k}{\frac{d^{2} m_{z}}{a^{2}-d^{2}}+m_{1}}
\end{gathered}
$$

\#2


$$
\begin{aligned}
& T=\frac{1}{2} m\left(\dot{x}^{2}+\dot{y}^{2}\right), V=m g y \\
& \quad \dot{y}=2 x \dot{x} \\
& T=\frac{1}{2} m\left(1+4 x^{2}\right) \dot{x}^{2}, V=m g x^{2}
\end{aligned}
$$

$$
\begin{aligned}
& \mathcal{L}=T-V=\frac{1}{2} m\left(1+4 x^{2}\right) \dot{x}^{2}-m g x^{2} \\
& \mathcal{L}=\frac{1}{2} f(x) \dot{x}^{2}-m g x^{2}, f(x)=\ln \left(1+4 x^{2}\right)
\end{aligned}
$$

\#2 continued
Similar to problem \#1,

$$
\frac{d}{d t} \frac{\partial l}{\partial \dot{x}}-\frac{\partial t}{\partial x}=0 \Rightarrow f \ddot{x}+f^{\prime} \frac{\dot{x}^{2}}{2}+2 m g x=0
$$

where $f(x)=m\left(1+4 x^{2}\right)$
linearize around $x=0$ :

$$
\begin{aligned}
& f(0) \ddot{x}+2 m g x=0, \quad f(0)=m \\
& \ddot{x}+2 g x=0 \\
& \omega^{2}=2 g, \quad \omega=\sqrt{2 g}
\end{aligned}
$$

\#3. $\quad x_{p}=R \cos (\Omega t-\phi)$

$$
\dot{x}_{p}=-R \Omega \sin (\Omega t-\phi)
$$

Amplitude of velocity $=R \Omega$
From class notes, $\frac{R}{\delta_{s t}}=\frac{1}{\sqrt{\left(1-\left(\frac{\left.\left.\delta_{0}^{2}\right)^{2}+4\left(\frac{n}{\omega}\right)\right)^{2}\left(\frac{2}{\omega}\right)^{2}}{}\right.\right.}}$

$$
\therefore \quad \frac{R \Omega}{\delta_{s t}}=\frac{\Omega}{\sqrt{\left(1-\left(\frac{s i}{\omega}\right)^{2}+4\left(\frac{n}{\omega}\right)^{2}\left(\frac{\pi}{w}\right)^{2}\right.}}
$$

diviansionles

$$
\begin{aligned}
\begin{array}{l}
\text { imensionles } \\
\text { ampitione of velocity }=\frac{R}{\delta_{s t}} \frac{\Omega}{\omega}
\end{array} & =\frac{\Omega / \omega}{\sqrt{\left(1-\left(\frac{s q^{2}}{w}\right)^{2}+4\left(\frac{n}{\omega}\right)^{2}\left(\frac{\Omega}{\omega}\right)^{2}\right.}} \\
& =\frac{q}{\sqrt{\left(1-q^{2}\right)^{2}+4 p^{2} q^{2}}}
\end{aligned}
$$

\# 3 continue d
Set $\frac{d}{d q}\left(\frac{q}{\sqrt{\left(1-q^{2}\right)^{2}+4 p^{2} q^{2}}}\right)=0$ for max
Obtain

$$
\frac{q^{4}-1}{\left(\left(1-q^{2}\right)^{2}+4 p^{2} q^{2}\right)^{3 / 2}}=0 \Rightarrow q=1
$$

$\therefore$ Max velocity amplitude occurs at

$$
\Omega=w
$$

\#4.

$$
\dot{\theta}=\sqrt{\frac{2 m g R}{m R^{2}+I_{0}}} \sqrt{\theta}, \theta(0)=0
$$

has two exact solutions:
Ore of than $i \quad \theta \equiv 0$ as stated.
The otha is obtained by sepal aton of variates

$$
\begin{aligned}
& \frac{d \theta}{\sqrt{\theta}}=\sqrt{\cdots} d t \\
& 2 \theta^{1 / 2}=\sqrt{\cdots} t+C^{7} \quad(c=0 \text { for } \theta(0)=0) \\
& \therefore \theta(t)=\frac{1}{2} \sqrt{\cdots} t^{2}
\end{aligned}
$$

\#4 continued
Comment: An equation with two (oumore) Solutions is bad news: how do you know if the solution you have found is the right one for your application?

That is why the mathematicians have given us a uniqueness proof which gives conditions on the right hand side of the differential equation $(f(0))$

$$
\frac{d \theta}{d t}=f(\theta)
$$

Which guarantee that the is only one solution.


$$
T=\frac{1}{2} m_{1} \dot{x}^{2}+\frac{1}{2} m_{2} V_{2}^{2}
$$

$$
\bar{r}_{2}=x \hat{e}_{x}+a \sin \theta \hat{e}_{x}
$$

$$
+a \cos \theta \hat{e}_{y}
$$

Let $m_{1}=M$

$$
\bar{V}_{2}=(\dot{x}+a \dot{\theta} \cos \theta,-a \dot{\theta} \sin \theta)
$$

$$
m_{2}=m
$$

$$
\bar{V}_{2}^{2}=(\dot{x}+a \dot{\theta} \cos \theta)^{2}+a^{2} \dot{\theta}^{2} \sin ^{2} \theta
$$

$$
=\dot{x}^{2}+2 a \dot{x} \dot{\theta} \cos \theta+a^{2} \dot{\theta}^{2}
$$

$$
V=\frac{1}{2} k x^{2}-a_{m g} \cos \theta
$$

$$
\begin{array}{r}
\frac{\partial L}{\partial x}=-k x \quad \frac{\partial k}{\partial \dot{x}}=m_{1} \dot{x}+m_{2}(\dot{x}+a \dot{\theta} \cos \theta) \\
\frac{d}{a t} "=\left[\begin{array}{c}
\left(m_{1}+m_{2}\right) \ddot{x}+m_{2}(a)\left(\ddot{\theta} \cos \theta-\dot{\theta}^{2} \sin \theta\right) \\
+k x=0
\end{array}\right.
\end{array}
$$

$$
\frac{\partial h}{\partial \theta}=-a m g \sin \theta-a_{x} \dot{\theta} \sin \theta
$$

$$
\frac{\partial h}{\partial \dot{\theta}}=\left(a \dot{x} \cos \theta+a^{2} \dot{\theta}\right) m_{2}
$$

$$
\frac{d}{d t} n=\left(a \ddot{x} \cos \theta-a \ddot{x} \quad \dot{\theta} \sin \theta+a^{2} \ddot{\theta}\right) m_{2}
$$

$$
m_{2}\left(a \ddot{x} \cos \theta+a^{2} \ddot{\theta}\right)+a m g \sin \theta=0
$$

$$
\begin{gathered}
(m+M) \ddot{x}+m^{a} \ddot{\theta} \quad+k x=0 \\
\ddot{x}+a \ddot{\theta}+g \theta=0
\end{gathered}
$$

(b) Linearize abut $x=\theta=0 \quad \cos \theta \rightarrow 1 \quad \sin \theta \rightarrow \theta$

1) $(m+m) \ddot{x}+m a \ddot{\theta}+k x=0$
2) $m a \ddot{x}+m a^{2} \ddot{\theta}+m g a \theta=0$
(C)

$$
\begin{array}{ll}
m \ddot{z}+K z=0 & z=\left[\begin{array}{l}
x \\
\theta
\end{array}\right] \\
\ddot{z}=\left[\begin{array}{l}
\ddot{x} \\
\ddot{\theta}
\end{array}\right] \\
M=\left[\begin{array}{cc}
(M+m) & m a \\
m a & m a^{2}
\end{array}\right] \quad K=\left[\begin{array}{cc}
K & 0 \\
0 & m g a
\end{array}\right]
\end{array}
$$

(d) $m=\left[\begin{array}{ll}2 & 1 \\ 1 & 1\end{array}\right]$
frequencies $\omega, \omega_{2}$

$$
\begin{aligned}
& \operatorname{det}\left[K-w^{2} M\right]=0 \\
& \operatorname{det}\left[\begin{array}{cc}
1-2 w^{2} & -w^{2} \\
-w^{2} & 1-w^{2}
\end{array}\right]=0 \\
& \left(1-2 w^{2}\right)\left(1-w^{2}\right)-w^{4}=0 \\
& 1-3 w^{2}+2 w^{4}-w^{4}=0 \\
& \left(w^{2}\right)^{2}-3\left(w^{2}\right)+1=0 \\
& w^{2}=\frac{3 \pm \sqrt{9-4}}{2} \\
& w_{1}^{2}=\frac{3-\sqrt{5}}{2} \quad w_{1}=0.6180 \\
& w_{2}^{2}=\frac{3+\sqrt{5}}{2} w_{2}=1.6180
\end{aligned}
$$

Modal vectors $i_{2}^{z}$

$$
m^{-1} K z=\omega^{2} z
$$

Solve by hand or in MATLAB

$$
1 z=\left[\begin{array}{l}
-0.5257 \\
-0.3249
\end{array}\right] \quad \text { a } z=\left[\begin{array}{c}
-0.85077 \\
1.3764
\end{array}\right]
$$

MATLAB: $[v, d]=\operatorname{eig}(K, M)$

$$
V=\left[\begin{array}{ll}
z & z^{z}
\end{array}\right]
$$

$$
d=\left[\begin{array}{cc}
w_{1}^{2} & 0 \\
0 & w_{2}^{2}
\end{array}\right]
$$

(e)

$$
\begin{aligned}
& \frac{\text { Orthogonality, condition }}{X_{j}^{\top} M X_{i}=0} \\
& X_{i}^{\top} K X_{i}=0
\end{aligned} \quad X_{1}=\left[\begin{array}{ll}
z & , z
\end{array}\right] \quad X_{2}=\left[\begin{array}{ll}
2 & 2
\end{array}\right]
$$

solving in MATLAB, all combinations of above equations yield zero.
(f)

$$
\begin{aligned}
& R=\left[\begin{array}{ll}
\exists & 2 \\
1 & z
\end{array}\right] \\
& R^{\top} M R=\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]
\end{aligned}
$$

diagonal
(9) $x, \theta$ as linear combinations of principal coordinates $p_{1}, p_{2}$ principal coordinates $p=p^{-1} z \quad(p g \mid 81)$

$$
\text { modal matrix } P=\left[\begin{array}{ll}
z_{2} & z
\end{array}\right] \quad z=P_{p} \quad z=\left[\begin{array}{l}
x \\
\theta
\end{array}\right]
$$

$$
\begin{aligned}
& x=-0.5257 p_{1}-0.8507 p_{2} \\
& \theta=-0.3249 p_{1}+1.3764 p_{2}
\end{aligned}
$$

b)

$$
\begin{aligned}
& R^{+} M R \ddot{p}+R^{\top} K R p=0 \\
& {\left[\begin{array}{cc}
1 & 0 \\
0 & 1
\end{array}\right] \ddot{p}+\left[\begin{array}{cc}
0.3820 & 0 \\
0 & 2.6180
\end{array}\right] p=0}
\end{aligned}
$$

(2) Neglect gravity

* 2 approaches

1) Model changing angles in crating EOM
2) Assume small vibrations in generation of EOM because M, K matrices will be linearized about small vibrations in calculation of $\omega_{1}, \omega_{2}$.

Approach 2


| component | $T$ RE | $V P E$ |
| :---: | :---: | :---: |
| $m$ | $\frac{1}{2} m\left(\dot{x}^{2}+\dot{y}^{2}\right)$ | - |
| $3 k$ | - | $\frac{1}{2} 3 k x^{2}$ |
| $2 k$ |  | $\frac{1}{2} 2 k\left(\Delta l_{2 k}\right)^{2}$ |
| $k$ |  | $\frac{1}{2} k\left(\Delta l_{k}\right)^{2}$ |

compression of $k$

$$
\left.\begin{aligned}
& \Delta l=[\cos 60, \sin 60] \cdot[x, y] \\
& D l=x \cos \frac{\pi}{3}+y \sin \frac{\pi}{3} \\
& \frac{O m p r s s i o n}{2} \text { of ak } \\
& D l_{2 k}=x \cos \frac{\pi}{3}-y \sin \frac{\pi}{3} \\
& 0.5 \\
& \sqrt{3} / 2
\end{aligned} \right\rvert\, \frac{d}{d t}
$$

$$
\begin{aligned}
& L=\frac{1}{2} m\left(\dot{x}^{2}+\dot{y}^{2}\right)-\frac{3}{2} k x^{2}-k \Delta l_{\partial t^{2}}^{2}-\frac{1}{2} k \Delta l_{k}^{2} \\
& \frac{d}{d t}\left(\frac{\partial L}{\partial \dot{x}}\right)-\frac{\partial L}{\partial x}=0 \quad \text { (1) } \\
& \frac{d}{d t}\left(\frac{\partial L}{\partial \dot{y}}\right)-\frac{\partial L}{\partial y}=0 \quad \text { (2) }
\end{aligned}
$$

$$
\Delta l_{k}{ }^{2}=\frac{1}{4} x^{2}+\frac{\sqrt{3}}{2} x y+\frac{3}{4} y^{2}
$$

$$
\Delta l_{2 k^{2}}^{2}=\frac{1}{4} x^{2}-\frac{\sqrt{3}}{2} x y+\frac{3}{4} y^{2}
$$

$$
\begin{aligned}
& \text { (1) } m \ddot{x}+3 k x+k\left(\frac{1}{2} x-\frac{\sqrt{3}}{2} y\right)+\frac{1}{2} k\left(\frac{1}{2} x+\frac{\sqrt{3}}{2} y\right)=0 \\
& m \ddot{x}+x\left(3 k+\frac{1}{2} k+\frac{1}{4} k\right)+y\left(-\frac{\sqrt{3}}{2} k+\frac{\sqrt{3}}{2} \cdot \frac{1}{2} k\right) \\
& \rightarrow m \ddot{x}+\frac{15}{4} k x-\frac{\sqrt{3}}{4} k y=0 \\
& \text { (2) } m \ddot{y}+k\left(-\frac{\sqrt{3}}{2} x+\frac{3}{2} y\right)+\frac{1}{2} k\left(\frac{\sqrt{3}}{2} x+\frac{3}{2} y\right)=0 \\
& \longrightarrow m \ddot{y}-\frac{\sqrt{3}}{4} k x+\frac{9}{4} k y=0 \\
& M=\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right] \quad K=\left[\begin{array}{cc}
15 / 4 & -\sqrt{3} / 4 \\
-\sqrt{3} / 4 & 9 / 4
\end{array}\right] \quad(\text { setting } K=1) \\
& \text { (Settion } m=1 \text { ) } \\
& \text { MATLAB }[v, d]=\operatorname{erg}(K, M) \\
& d=\left[\begin{array}{cc}
w_{1}{ }^{2} & 0 \\
0 & w_{2}^{2}
\end{array}\right] \\
& \omega_{1}=1.4608 \\
& \omega_{2}=1.9662
\end{aligned}
$$

2. 


$T=\frac{1}{2} m\left(\dot{x}^{2}+\dot{y}^{2}\right)$
Lagrange's Es give

$$
\begin{aligned}
& m \ddot{x}=-\frac{\partial v}{\partial x} \\
& m \ddot{y}=-\frac{\partial v}{\partial y}
\end{aligned}
$$

Blow up view:


$$
\begin{aligned}
V= & \frac{1}{2} k\left(\sqrt{\left(\frac{1}{2}-x\right)^{2}+\left(\frac{\sqrt{3}}{2}-y\right)^{2}}-1\right)^{2} \\
& +\frac{1}{2}(2 k)\left(\sqrt{\left(\frac{1}{2}-x\right)^{2}+\left(\frac{\sqrt{3}}{2}+y\right)^{2}}-1\right)^{2} \\
& +\frac{1}{2}(3 k)\left(\sqrt{(1+x)^{2}+y^{2}}-1\right)^{2}
\end{aligned}
$$

Expand $V$ using the Taylor series for $\sqrt{1+Z}$

$$
\begin{aligned}
& \sqrt{1+z}= 1+\frac{z}{2}-\frac{z^{2}}{8}+\cdots \\
& \sqrt{\left(\frac{1}{2}-x\right)^{2}+\left(\frac{\sqrt{3}}{2}-y\right)^{2}}= \sqrt{\frac{1}{4}-x+x^{2}+\frac{3}{4}-\sqrt{3} y+y^{2}} \\
&= \sqrt{1+\left(-x+x^{2}-\sqrt{3} y+y^{2}\right)} \\
&= 1+\frac{\left(-x+x^{2}-\sqrt{3} y+y^{2}\right)}{2} \\
&-\frac{\left(-x+x^{2}-\sqrt{3} y+y^{2}\right)^{2}}{8}+\cdots \\
&=\left.1-\frac{1}{8}-\frac{\sqrt{3} y}{2}+\frac{x^{2}}{2}+\frac{y^{2}}{2}+3 y^{2}+2 \sqrt{3} x y\right)+/ \\
&= 1-\frac{x}{2}-\frac{\sqrt{3}}{2} y+\frac{3}{8} x^{2}+\frac{1}{8} y^{2}-\frac{\sqrt{3}}{4} \times y \\
&= \frac{1}{4}-x+x^{2}+\frac{3}{4}-\sqrt{3} y+y^{2}+1 \\
&\left(\sqrt{\left.\left(\frac{1}{2}-x\right)^{2}+\left(\frac{\sqrt{3}}{2}-y\right)^{2}-1\right)^{2}}=\right.-2+x+\sqrt{3} y-\frac{3}{4} x^{2}-\frac{1}{4} y^{2}+\frac{\sqrt{3}}{2} x y \\
&=\left.x^{2}+y^{2}-\frac{1}{4}-y\right)^{2}+1-2\left(\frac{1}{4} x^{2}+\frac{3}{4} y^{2}+\frac{\sqrt{3}}{2} x y\right. \\
&= \frac{\sqrt{3}}{2} x y \\
&=
\end{aligned}
$$

$$
\begin{aligned}
& \sqrt{(1+x)^{2}+y^{2}}=\sqrt{1+2 x+x^{2}+y^{2}} \\
&=1+\frac{1}{2}\left(2 x+x^{2}+y^{2}\right)-\frac{1}{8}\left(2 x+x^{2}+y^{2}\right)^{2}+\cdots \\
&=1+x+\frac{x^{2}}{2}+\frac{y^{2}}{2}-\frac{x^{2}}{2} \\
&=1+x+\frac{y^{2}}{2} \\
&\left.\left.\begin{array}{rl}
\left.\sqrt{(1+x)^{2}+y^{2}}-1\right)^{2} & =(1+x)^{2}+y^{2}+1-2\left(1+x+\frac{y^{2}}{2}\right) \\
& =1+2 x+x^{2}+y^{2}+1-2-2 x-y^{2} \\
& =x^{2} \\
V=\frac{1}{2} k\left(\frac{x^{2}}{4}+\frac{3}{4} y^{2}+\frac{\sqrt{3}}{2} x y\right) \\
+\frac{1}{2}(2 k)\left(\frac{x^{2}}{4}\right. & \left.+\frac{3}{4} y^{2}-\frac{\sqrt{3}}{2} x y\right) \\
V=\frac{1}{2} k[3 k) x^{2} \\
=\frac{x^{2}}{2}\left(\frac{1}{4}\right. & \left.+\frac{2}{4}+3\right)+\frac{15}{4} x^{2}
\end{array}\right)+\frac{9}{4} y^{2}-\frac{\sqrt{3}}{2} x y\right] \\
& m \ddot{x}=-\frac{2 v}{2 x}=-\frac{15}{4} k x+\frac{\sqrt{3}}{4} y k \\
& m=-\frac{\partial v}{2 y}=-\frac{9}{4} k y+\frac{\sqrt{3}}{4} x k \\
&\left.=\sqrt{3} x y\left(\frac{1}{2}-1\right)\right]
\end{aligned}
$$

$$
\left.\begin{gathered}
M \ddot{x}+K x=0 \\
{\left[\begin{array}{cc}
m & 0 \\
0 & m
\end{array}\right]\left[\begin{array}{c}
\ddot{x} \\
\ddot{y}
\end{array}\right]+\left[\begin{array}{cc}
\frac{15}{4} k & -\frac{\sqrt{3}}{4} k \\
-\frac{\sqrt{3}}{4} k & \frac{9}{4} k
\end{array}\right]\left[\begin{array}{l}
x \\
y
\end{array}\right]=\left[\begin{array}{l}
0 \\
0
\end{array}\right]} \\
x=A \cos \omega t \\
y=B \cos \omega t \\
\left\lvert\,-\omega^{2} m+\frac{15}{4} k \quad-\frac{\sqrt{3}}{4} k\right. \\
-\frac{\sqrt{3}}{4} k \quad-\omega^{2} m+\frac{9}{4} k
\end{gathered} \right\rvert\,=0 \quad \begin{aligned}
& \left.\beta^{2}+\frac{15}{4}\right)\left(-\beta^{2}+\frac{9}{4}\right)-\frac{3}{16}=0 \quad \omega h e r e \beta=\omega \sqrt{\frac{m}{k}} \\
& \beta^{4}-6 \beta^{2}+\frac{33}{4}=0 \\
& \beta^{2}=\frac{6 \pm \sqrt{36-33}}{2}=3 \pm \frac{\sqrt{3}}{2} \\
& \omega=\sqrt{3 \pm \frac{\sqrt{3}}{2}} \sqrt{\frac{k}{m}}=1.461 \sqrt{\frac{k}{m}}, 1.966 \sqrt{\frac{k}{m}}
\end{aligned}
$$

3. The general motion of the first coordinate of a two degree of freedom system is given by:

$$
x_{1}(t)=R_{1} \cos \left(\omega_{1} t-\theta_{1}\right)+R_{2} \cos \left(\omega_{2} t-\theta_{2}\right)
$$

Is this a periodic motion? Under what condition will it be periodic?
At $t=0$,

$$
x_{1}(0)=R_{1} \cos \left(\theta_{1}\right)+R_{2} \cos \left(\theta_{2}\right)
$$

At what time $t$ will this happen again?
Suppose that $\omega_{2}=\frac{m}{n} \omega_{1}$, where $m$ and $n$ are whole numbers. Then

$$
x_{1}(t)=R_{1} \cos \left(\omega_{1} t-\theta_{1}\right)+R_{2} \cos \left(\frac{m}{n} \omega_{1} t-\theta_{2}\right)
$$

After time $T=\frac{2 \pi n}{\omega_{1}}$, we have

$$
x_{1}(T)=R_{1} \cos \left(2 \pi n-\theta_{1}\right)+R_{2} \cos \left(2 \pi m-\theta_{2}\right)=x_{1}(0)
$$

In fact, $x_{1}(T+t)=x_{1}(t)$ for all $t$, not just $t=0$. Thus in this case the motion is periodic.
However, if the ratio of $\omega_{2}$ to $\omega_{1}$ is an irrational number, then $x_{1}(t)$ will never return to $x_{1}(0)$ and the motion will not be periodic.
(1)


$$
\begin{aligned}
& T=\frac{1}{2}\left(\dot{x}_{1}^{2}+\dot{x}_{2}^{2}\right) \\
& V=\frac{1}{2}\left(x_{1}^{2}+x_{2}^{2}+\left(x_{1}-x_{2}\right)^{2}\right) \\
& =x_{1}^{2}+x_{2}^{2}-x_{1} x_{2} \\
& \delta W_{2}=F \cos \Omega t \delta x_{2} \Rightarrow Q_{2}=F \cos \Omega t \\
& \delta W_{1}=0 \Rightarrow Q_{1}=0 \\
& \ddot{x}_{1}+2 x_{1}-x_{2}=0 \\
& \ddot{x}_{2}-x_{1}+2 x_{2}=F \cos \Omega t \\
& M=\left(\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right), K=\left(\begin{array}{cc}
2 & -1 \\
-1 & 2
\end{array}\right), f=\binom{0}{F \cos \Omega t} \\
& M \ddot{x}+K x=f(t) \\
& \text { Let } x=\Sigma \cos \omega t \text { for } f(t)=0 \\
& (-\omega \bar{I}+k) Z=0 \\
& \left|\begin{array}{cc}
-w^{2}+2 & -1 \\
-1 & -w^{2}+2
\end{array}\right|=0,\left(-w^{2}+2\right)^{2}=1 \\
& -\omega^{2}+2= \pm 1 \\
& \omega^{2}=2 \mp 1=3,1 \\
& \omega_{1}=1, \quad\left[\begin{array}{cc}
-1+2 & -1 \\
-1 & -1+2
\end{array}\right], \bar{X}=0 \Rightarrow, \bar{X}=\binom{1}{1} \\
& w_{2}=\sqrt{3}, \quad\left[\begin{array}{ll}
-1 & -1 \\
-1 & -1
\end{array}\right], \mathbb{Z}=0 \Rightarrow \underset{2}{Z}=\binom{1}{-1}
\end{aligned}
$$

Set $x=R p, \quad R=[1 \Phi, 2 X]=\left[\begin{array}{rr}1 & 1 \\ 1 & -1\end{array}\right]$

$$
\begin{gathered}
x_{1}=p_{1}+p_{2} \\
x_{2}=p_{1}-p_{2} \\
\underbrace{R^{t} M R}_{R^{*} I R} \ddot{p}+\underbrace{\left[\begin{array}{rr}
1 & 1 \\
1 & -1
\end{array}\right]}_{\left[\begin{array}{ll}
2 & 0 \\
0 & 2
\end{array}\right]} \begin{array}{rr}
1 & 1 \\
1 & -1
\end{array}]
\end{gathered}+\left[\begin{array}{rr}
1 & 1 \\
1 & -1
\end{array}\right] \underbrace{\left[\begin{array}{cc}
2 & -1 \\
-1 & 2
\end{array}\right]\left[\begin{array}{cc}
1 & 1 \\
1 & -1
\end{array}\right]=\left[\begin{array}{cc}
2 & 0 \\
0 & 6
\end{array}\right]}_{\left[\begin{array}{cc}
1 & 3 \\
1 & -3
\end{array}\right]}=
$$

$$
\begin{aligned}
& \ddot{p}_{1}+\omega_{1}^{2} p_{1}=\frac{F}{2} \cos \Omega t, \quad \omega_{1}=1 \\
& \ddot{p}_{2}+\omega_{2}^{2} p_{2}=-\frac{F}{2} \cos \Omega t, \quad \omega_{2}=\sqrt{3} \\
& p_{1}=k \cos \Omega t, \quad\left(-\Omega^{2}+\omega_{1}^{2}\right) k=F
\end{aligned}
$$

So $\quad p_{1}=\frac{F / 2}{1-\Omega^{2}} \cos \Omega t$
Similasly $p_{2}=\frac{-F / 2}{3-\Omega^{2}} \cos \Omega t$

$$
\begin{aligned}
\therefore x_{1}=p_{1}+p_{2} & =\left(\frac{1}{1-\Omega^{2}}-\frac{1}{3-\Omega^{2}}\right) \frac{F}{2} \cos \Omega t \\
& =\frac{F}{\left(1-\Omega^{2}\right)\left(3-\Omega^{2}\right)} \cos \Omega t \\
x_{2}=p_{1}-p_{2} & =\left(\frac{1}{1-\Omega^{2}}+\frac{1}{3-\Omega^{2}}\right) \frac{F}{2} \cos \Omega t \\
& =\frac{\left(2-\Omega^{2}\right) F}{\left(1-\Omega^{2}\right)\left(3-\Omega^{2}\right)} \cos \Omega t
\end{aligned}
$$

$$
\begin{aligned}
& \left.\begin{array}{l}
\ddot{x}_{1}+2 x_{1}-x_{2}=0 \\
\ddot{x}_{2}-x_{1}+2 x_{2}=F \cos \Omega t \\
\text { Set } \begin{array}{l}
x_{1}
\end{array}=A \cos \Omega t \\
x_{2}=B \cos \Omega t \\
-\Omega^{2} A+2 A-B=0 \\
-\Omega^{2} B-A+2 B=F
\end{array}\right\} \text { Solve for } A, B \\
& \Rightarrow \quad A=\frac{F}{\left(1-\Omega^{2}\right)\left(3-\Omega^{2}\right)}, B=\frac{F\left(2-\Omega^{2}\right)}{\left(1-\Omega^{2}\right)\left(3-\Omega^{2}\right)}
\end{aligned}
$$

Agrees with (1)

SOLUTION to question 3:
Multiply the first eq. in (7) by $-\Omega^{2}$ and add to the second eq. in (7) giving:

$$
R^{t}\left(-\Omega^{2} M+K\right) R=-\Omega^{2} D_{1}+D_{2}
$$

Take the inverse of both sides:

$$
\begin{gathered}
\left(R^{t}\left(-\Omega^{2} M+K\right) R\right)^{-1}=\left(-\Omega^{2} D_{1}+D_{2}\right)^{-1} \\
R^{-1}\left(-\Omega^{2} M+K\right)^{-1}\left(R^{t}\right)^{-1}=\left(-\Omega^{2} D_{1}+D_{2}\right)^{-1}
\end{gathered}
$$

Now multiply on the left by $R$ and on the right by $R^{t}$, giving

$$
\left(-\Omega^{2} M+K\right)^{-1}=R\left(-\Omega^{2} D_{1}+D_{2}\right)^{-1} R^{t}
$$

This demonstrates the equivalence of eqs.(5) and (13).


$$
u_{t t}=c^{2} u_{x x}, u_{x}=0 \text { at } x=0, l
$$

Set $u=U(x) \cos \omega t$

$$
\begin{aligned}
-\omega^{2} U & =c^{2} U^{\prime \prime} \\
U(x) & =c_{1} \sin \frac{\omega}{c} x+c_{2} \cos \frac{\omega}{c} x \\
U^{\prime} & =\frac{\omega}{c}\left(c_{1} \cos \frac{\omega}{c} x-c_{2} \sin \frac{\omega}{c} x\right) \\
U^{\prime}(0) & =U^{\prime}(l)=0 \Rightarrow c_{1}=0 \text { and }
\end{aligned}
$$

$\sin \frac{\omega l}{c}=0$

$$
\frac{\omega l}{c}=n \pi, n=0,1,2, \ldots
$$

a)

$$
\omega_{n}=n \pi \frac{c}{l}, U_{n}(x)=\cos \frac{\omega_{n}}{c} x
$$

b) Show $\left\{V_{n}\right\}$ is or the goral

$$
\begin{aligned}
& \int_{0}^{l} U_{n} U_{m} d x=0, n \neq m \\
& \int_{0}^{l} \underbrace{\cos \frac{n \pi x}{l} \cos \frac{m \pi x}{l}} d x=0 \\
& \frac{1}{2} \cos \left(\left(\frac{n+m}{I}\right)^{\pi} x\right)+\frac{1}{2} \cos \left(\left(\frac{n-m}{l}\right)^{\pi} x\right) \\
& =\frac{l}{2 \pi}\left(\frac{\sin \left(\frac{n+m}{x} x\right)}{n+m}+\left.\frac{\sin \left(\frac{n-m}{1, x}\right)}{n-m}\right|_{0} ^{l}\right) \\
& =\frac{\sin (n+m) \pi}{n+m}+\frac{\sin (n-m) \pi}{n-m}=0
\end{aligned}
$$

C)

$$
u(x, t)=\left\{\begin{array}{c}
\sum_{n=1}^{\infty}\left(a_{n} \cos \omega_{n} t+b_{n} \sin \omega_{n} t\right) \cos \frac{\omega_{n} x}{c} x \\
+a_{0}+b_{0} t
\end{array}\right.
$$

Trigid body mode
d) IC $t=0, u_{t}=0 \Rightarrow b_{n}=0, n=0,1,2, \ldots$

$$
t=0, \quad u=\frac{x}{l}=a_{0}+\sum_{n=1}^{\infty} a_{n} \cos \frac{n \pi x}{l}
$$

Mull.by $\cos \frac{m \pi x}{l} \& \int_{0}^{l} \Rightarrow$

$$
\begin{aligned}
\int_{0}^{l} \frac{x}{l} \cos \frac{m \pi x}{l} d x=\int_{0}^{l} a_{m}\left(\cos \frac{m \pi x}{l}\right)^{2} d x=\frac{l}{2} a_{m} \\
(m>0)
\end{aligned} a_{m}=\frac{2}{l} \int_{0}^{l} \frac{x}{l} \cos \frac{m \pi x}{l} d x, m>0 . ~\left(\frac{1}{l} \int_{0}^{l} \frac{x}{l} d x=\frac{1}{l^{2}}\left[\left.\frac{x^{2}}{2}\right|_{0} ^{l}\right]=\frac{1}{2} .\right.
$$

e)

$$
\begin{aligned}
& \text { at } x=\frac{l}{2}, u\left(\frac{l}{2}, t\right)=\frac{1}{2}+\sum_{n=1,3,5}^{\infty} a_{n} \cos \omega_{n} t \cos \frac{n \pi}{2}=\frac{1}{2} \\
& \text { at } x=0, u(0, t)
\end{aligned}=\frac{1}{2}+\sum_{n=1,2,5}^{\infty} a_{n} \cos \frac{n \pi c t}{l}, \begin{aligned}
\text { at } x=l, u(l, t) & =\frac{1}{2}+\sum_{n=1,3,5}^{\infty} a_{n} \cos \frac{n \pi c t}{2} \cos n n^{-1} \\
& =\frac{1}{2}-\sum_{n=1,1,5,5}^{\infty} a_{n} \cos \frac{n \pi t t}{2}
\end{aligned}
$$

f)

$$
\begin{array}{r}
x(0, t)=\frac{1}{2}-\frac{4}{\pi^{2}}\left(\cos \frac{\pi c}{l} t+\frac{1}{9} \cos \frac{3 \pi l}{l} t\right. \\
\\
\left.+\frac{1}{25} \cos \frac{5 \pi \pi^{2}}{l} t+\cdots\right)
\end{array}
$$

See figine attriched.
2.


See text, prottem 7.8, pp. 210-211
3.


$$
E A \frac{\partial u}{\partial x}=-M \frac{\partial^{2} u}{\partial t^{2}} \text { at } x=l
$$

seetexts problem 7.4, p. 207

(a) $\omega_{n}=n \pi, \quad U_{n}(x)=\cos n \pi x$
b)

$$
\begin{aligned}
f(x, t) & =x^{2} \cos t=\sum f_{i}(t) U_{i}(x) \\
f_{n}(t) & =\left(\frac{\int_{0}^{1}(\cos n \pi x) x^{2} d x}{\int_{0}^{1}(\cos n \pi x)^{2} d x}\right) \cos t \\
& =\frac{\frac{2}{\pi^{2} n^{2}}(-1)^{n}}{\frac{1}{2}} \cos t=F_{n} \cos t
\end{aligned}
$$

where $F_{n}=\frac{4}{\pi^{2} n^{2}}(-1)^{n}, n>0 . \quad F_{0}=\frac{1}{3}$
c)

$$
\begin{aligned}
& H(x)= u_{t}(x, 0)=x=\sum H_{n} \cos n \pi x \\
& H_{n}=\frac{\int_{0}^{1} x \cos n \pi x d x}{\int_{0}^{1}(\cos n \pi x)^{2} d x}=\frac{\frac{1}{\pi^{2} n^{2}\left((-1)^{n}-1\right)}}{\frac{1}{2}}, n>0 \\
& H_{n}=\left\{\begin{array}{c}
-\frac{4}{\pi^{2} n^{2}}, n \cdot d d \\
0, n \text { even }
\end{array}\right\} n>0 ; H_{0}=\int_{0}^{1} x d x=\frac{1}{2} \\
& G(x)=n(x, 0)=0 \Rightarrow G_{n}=0
\end{aligned}
$$

d) $\quad p_{n}^{\prime \prime}+w_{n}^{2} p_{n}=F_{n} \cos t$
$t=0, p_{n}=0, \dot{p}_{n}=H_{n}$
e) $\quad p_{n}(t)=A_{n} \cos n \pi t+B_{n} \sin n \pi t+\frac{F_{n} \cos t}{(n \pi)^{2}-1}, n>0$
(contisined) $\quad p_{0}(t)=\frac{1}{2} t-\frac{1}{3} \cos t+\frac{1}{3}$
where $A_{n}=\frac{-F_{n}}{(n \pi)^{2}-1}$

$$
B_{n}=\frac{H_{n}}{n \pi}
$$

where $F_{n}$ and $H_{n}$ are given in b) and c)
f)

$$
\begin{aligned}
& u(x, t) \approx p_{1}(t) U_{1}(x)+p_{0}(t) U_{0}(x) \\
& \approx\left(A_{1} \cos \pi t+B_{1} \sin \pi t+\frac{F_{1} \cos t}{\pi^{2}-1}\right) \cos \pi x+p_{0}(t) \\
& \approx\left(\frac{\left.-\frac{4}{\pi^{2}\left(\pi^{2}-1\right)}(\cos t-\cos \pi t)-\frac{4}{\pi^{3}} \sin \pi t\right) \cos \pi x}{} \quad+\frac{1}{2} t-\frac{1}{3} \cos t+\frac{1}{3}\right. \\
&2 a) \quad u=U(x) \cos \omega t \\
&-\omega^{2} U+\frac{E I}{5} U^{N}=0 \\
& U^{\prime v}-k^{4} U=0, \quad k^{4}=\omega^{2} \frac{\rho}{E I} \\
& U=e^{\lambda x} \Rightarrow \quad \lambda^{4}-k^{4}=0 \quad \Rightarrow \lambda=k,-k, i k,-i k \\
& U=C_{1} \cosh k x+c_{2} \sinh k x+c_{3} \cos k x+c_{4} \sin k x \\
& B C \quad U(0)=U(l)=U^{\prime}(0)=U^{\prime}(l)=0
\end{aligned}
$$

4 homegeneans algebraic ens.
For nontrivial solution, set determinant $=0$ which gives

$$
\begin{aligned}
& \cos k l \cosh k l=1, \quad k=\sqrt{\omega}\left(\frac{\rho}{E I}\right)^{1 / 4} \\
& \omega=\frac{(k l)^{2}}{l^{2}}\left(\frac{E I}{5}\right)^{1 / 2}
\end{aligned}
$$

Solving the system of 4 hanog. es, obtain

$$
\begin{aligned}
U_{n}(x)= & \cosh (k l) \frac{x}{l}-\cos (k l) \frac{x}{l} \\
& +\mu\left(\sinh \left(k_{l}\right) \frac{x}{l}-\sin (k l) \frac{x}{l}\right)
\end{aligned}
$$

where

$$
\mu=-\frac{(\cosh k \ell-\cos h l)}{(\sinh k \ell-\sin h l)}
$$

b) $\quad k l=4.73,7.85,10.99,14.13$
c) See plot, attached.

Note: This information may be obtained directly from the "Table of Beam Frequencies" posted on the web.

At the bottom of the Table we fried:

$$
\begin{aligned}
& \omega=\frac{\lambda^{2}}{\ell^{2}} \sqrt{\frac{E I}{\rho}} \text { where } \lambda=k \ell \text { in above notation } \\
& J(u)=\cosh u-\cos u \\
& H(x)=\sinh u-\sin u
\end{aligned}
$$

Using this notation, the Tall gives the mode shape as

$$
U_{n}(x)=J\left(\lambda_{n} \frac{x}{l}\right)-\frac{J\left(\lambda_{n}\right)}{H\left(\lambda_{n}\right)} H\left(\lambda_{n} \frac{x}{l}\right)
$$


3.

$$
\begin{aligned}
& \frac{d^{2} u}{d x^{2}}+u=1 \\
& u(0)=0 \\
& u(\pi)=0
\end{aligned}
$$

General solution

$$
\begin{aligned}
u & =A \sin x+B \cos x+1 \\
u(0) & =B+1=0 \Rightarrow B=-1 \\
u(\pi) & =-B+1=0 \Rightarrow B=1
\end{aligned}
$$

Since B cannot be equal to both 1 and -1 , this problem has no solution.

Note that the above syptem does have a solution for appropilite choices of the Jugitt hard side. For example

$$
\left.\left.\begin{array}{rl}
\frac{d^{2} u}{d x^{2}}+u=\cos 3 x \\
u & =A \sin x+B \cos x-\frac{1}{8} \cos 3 x \\
u(0) & =B-\frac{1}{8}=0 \Rightarrow B=1 / 8 \\
u(\pi) & =-B-\frac{1}{8}(-1)=0 \Rightarrow B=1 / 8
\end{array}\right\} \begin{array}{l}
\text { no }
\end{array}\right\}
$$

This can be exp pained in terms of the "Fredholm alternative theorem" which I will go over in class.
la.


From Table, $\lambda_{1}=3.9266$

$$
\omega_{1}=\lambda_{1}^{2} \sqrt{\frac{E I}{\rho l^{4}}}=15.42 \sqrt{\frac{E I}{\rho l^{4}}}
$$

16. 

$$
\begin{aligned}
& \frac{d^{4} u}{d x^{4}}=1 \\
& u=c_{1}+c_{2} x+c_{3} x^{2}+c_{4} x^{3}+\frac{x^{4}}{24}
\end{aligned}
$$

$B C$ :

$$
\begin{aligned}
& x=0, u=0, u^{\prime}=0 \Rightarrow c_{1}=c_{2}=0 \\
& x=l, u=0, u^{\prime \prime}=0 \Rightarrow c_{3}=\frac{l^{2}}{16}, c_{4}=-\frac{5}{48} l \\
& u=\frac{x^{2} l^{2}}{16}-\frac{5}{48} x^{3} l+\frac{x^{4}}{24}(=V(x) \text { below })
\end{aligned}
$$

ic.

$$
\begin{aligned}
& Q=\frac{E I \int_{0}^{l}\left(V^{\prime \prime}\right)^{2} d x}{\rho \int_{0}^{l} V^{2} d x}=\frac{E I \frac{l^{5}}{320}}{\rho \frac{19 l^{9}}{1451520}} \\
& \omega_{1}<\sqrt{Q}=\sqrt{\frac{4536}{19}} \sqrt{\frac{E I}{\rho l^{4}}}=15.45 \sqrt{\frac{E I}{\rho l^{4}}}
\end{aligned}
$$

Great agreement with $1 a$ !

1d. Again take $V(x)=\frac{x^{2} l^{2}}{16}-\frac{5}{48} x^{3} l+\frac{x^{4}}{24}$


$$
\begin{aligned}
& \rho=\rho_{0}+m \delta\left(x-\frac{l}{2}\right) \\
& Q=\frac{E I \int_{0}^{l}\left(V^{\prime \prime}\right)^{2} d x}{\rho_{0} \int_{0}^{l} V^{2} d x+m V\left(\frac{l}{2}\right)^{2}} \\
& V\left(\frac{l}{2}\right)=\frac{l^{4}}{192} \\
& Q=\frac{E I \frac{l^{5}}{320}}{\rho_{0} \frac{19 l^{9}}{1451520}+\frac{m l^{8}}{(142)^{2}}} \\
&=\frac{E I\left(\frac{4536}{19}\right)}{\rho \cdot l^{4}+\left(\frac{315}{152}\right) m l^{3}} \\
& \omega_{1}<\sqrt{Q}=15.45 \sqrt{\frac{E I}{\rho \cdot l^{4}+2.07 m l^{3}}}
\end{aligned}
$$

2. $r(x)=x+1 \quad$ (See 7.22 on p.225)

$$
\begin{aligned}
& A(x)=\pi(x+1)^{2} \\
& Q=\frac{\int_{0}^{2} E A(x)\left(u^{\prime}\right)^{2} d x}{\int_{0}^{2} \rho A(x) u^{2} d x}
\end{aligned}
$$

Choose $u(x)=x(x-2)$
Which satisfies the $B C \quad u(0)=u(2)=0$
Then $u^{\prime}=2 x-2$

$$
\begin{aligned}
& Q=\frac{E \int_{0}^{2} \pi(x+1)^{2}(2 x-2)^{2} d x}{\rho \int_{0}^{2} \pi(x+1)^{2} x^{2}(x-2)^{2} d x}=\frac{\frac{184}{15} E}{\frac{464}{105} \rho} \\
& Q=\frac{161}{58} \frac{E}{5} \Rightarrow \omega_{1} \leqslant \sqrt{Q}=1.67 \sqrt{\frac{E}{5}}
\end{aligned}
$$

3. $u^{\prime v}-u=1$
a) $\quad u=c_{1} \sin x+c_{2} \cos x+c_{3} \sinh x+c_{4} \cosh x-1$

$$
\begin{aligned}
& \left.\begin{array}{l}
u(0)=0=c_{2}+c_{4}-1 \\
u^{\prime \prime}(0)=0=-c_{2}+c_{4}
\end{array}\right\} \quad c_{2}=c_{4}=\frac{1}{2} \\
& \left\{\begin{array}{l}
u(\pi)=0=c_{4} \tilde{c}+c_{3} \tilde{s}-c_{2}-1 \quad \text { where } \begin{array}{r}
\tau=\cosh (\pi) \\
\\
u^{\prime \prime}(\pi)=0=\sinh (\pi)
\end{array}
\end{array} . \begin{array}{l}
u c_{4} \tilde{c}+c_{3} \tilde{s}+c_{2}
\end{array}\right.
\end{aligned}
$$

Substituting $C_{2}=C_{4}=\frac{1}{2}$ from above gives 2 incompatible values for $C_{3} \Rightarrow$ no solution
b) The Fredholm Aftunative sago $f(x)$ must be outhegosal to the null space of the adjoint.
[Not ethat (i) $u^{\prime N}-u=0 \Rightarrow u=\sin x$ satisfies $B_{1} C$. ie. the homogeneous syptum has a nontrivial solo and lii] the opuator $L u=u^{\prime v} u$ is setf-adjoint.] Only those $f(x)$ munch satisfy $\int_{0}^{\pi} f(x) \sin x d x=0$ will give a solution, $E_{1 g}$. $f(x)=\cos 3 x$.

Consider a clamped-free beam of constant depth, and a width which varies linearly from a maximum at the fixed end to zero at the free end. Taking the origin of coordinates at the fixed end, $I=I_{0}(1-x / l)$ and $\mu=\mu_{0}(1-x / l)$, where $I_{0}$ and $\mu_{0}$ are respectively the moment of inertia and mass per unit length at the fixed end. Choose a two-term series,

$$
W=A_{1} x^{2}+A_{2} x^{3}
$$

The kinetic and potential energies are:

$$
\begin{aligned}
T^{*} & =32 \int_{0}^{l} \mu W^{2} d x=\frac{\mu_{0}}{2} \int_{0}^{l}\left(1-\frac{x}{l}\right)\left(A_{1} x^{2}+A_{2} x^{3}\right)^{2} d x \\
V_{\max } & =1 / 2 \int_{0}^{l} E I\left(W^{\prime \prime}\right)^{2} d x=\frac{E I_{0}}{2} \int_{0}^{l}\left(1-\frac{x}{l}\right)\left(2 A_{1}+6 A_{2} x\right)^{2} d x
\end{aligned}
$$

Integrating, taking the partial derivatives, and substituting into (61.127) gives the pair of equations,

$$
\begin{align*}
& (2-\beta / 30) A_{1}+(2-\beta / 42) A_{2} l=0  \tag{61.128a,b}\\
& (2-\beta / 42) A_{1}+(3-\beta / 56) A_{2} l=0
\end{align*}
$$

where $\beta=\mu_{0} l^{1} \lambda / E I_{0}$. The roots of the frequency equation are $\beta_{1}=51.25$ and $\beta_{2}=1377$, and, hence, the first two frequencies are $\omega_{1}^{2} \leq 51.25 E I_{0} / \mu_{0} l^{4}$ and $\omega_{2}^{2} \leq$ $1377 E I_{0} / \mu_{0} l^{4}$
2. Ritzon
w/ 3 tams

$$
\begin{aligned}
& V=c_{1} x^{2}+c_{2} x^{3}+c_{3} x^{4} \\
& \bar{Q}=Q \frac{\rho l^{4}}{E I}, \bar{c}_{2}=c_{2} l, \bar{c}_{3}=c_{3} l^{2} \\
& \bar{Q}=l^{4} \frac{\int_{0}^{l}\left(v^{\prime \prime}\right)^{2} d x}{\int_{0}^{l} v^{2} d x} \\
& \int_{0}^{l}\left(V^{\prime \prime}\right)^{2} d x=l\left(\frac{144}{5} \bar{c}_{3}^{2}+36 \bar{c}_{2} \bar{c}_{3}+16 c_{1} \bar{c}_{3}+12 \bar{c}_{2}^{2}\right. \\
& \left.+12 c_{1} \bar{c}_{2}+4 c_{1}^{2}\right)=l F \\
& \int_{0}^{l} v^{2} d x=l^{5}\left(\frac{\bar{c}_{3}^{2}}{9}+\frac{\bar{c}_{2} \bar{c}_{3}}{4}+\frac{2}{7} c_{1} \bar{c}_{3}+\frac{\bar{c}_{2}^{2}}{7}\right. \\
& \left.+\frac{c_{1} \bar{c}_{2}}{3}+\frac{c_{1}^{2}}{5}\right)=\ell^{5} G \\
& \bar{Q}=\frac{F}{G}, \quad G \bar{Q}=F
\end{aligned}
$$

take $\frac{2}{x_{1}}, \frac{2}{x_{2}} * \frac{2}{x_{3}}$ ot his eq. $\Leftrightarrow$ set $\frac{\partial \bar{Q}}{\partial C_{i}}=0$

$$
\left[\begin{array}{ccc}
\frac{2 \bar{Q}-40}{5} & \frac{\bar{Q}-36}{3} & \frac{2 \bar{Q}-112}{7} \\
\frac{\bar{Q}-36}{3} & \frac{2 \bar{Q}-168}{7} & \frac{\bar{Q}-144}{4} \\
\frac{2 \bar{Q}-112}{7} & \frac{\bar{Q}-144}{4} & \frac{10 \bar{Q}-2512}{45}
\end{array}\right]\left[\begin{array}{l}
C_{1} \\
C_{2} \\
C_{3}
\end{array}\right]=\overline{0}
$$

Set $\operatorname{det}=0 \Rightarrow$

$$
\begin{aligned}
& 5 \bar{Q}^{3}-72324 \bar{Q}^{2}+35392896 \bar{Q}-426746880=0 \\
& \bar{Q}=12.369,494.322,13958.107
\end{aligned}
$$

ExACT:

$$
\begin{array}{ll}
\bar{w}_{1} \leq \sqrt{\bar{Q}}=3.517 & \text { (vs. } \left.1.8851^{2}=3.516\right) \\
\bar{w}_{2} \leq \sqrt{\bar{Q}}=22.233 & \left(\text { vs. } 4.6441^{2}=22.03\right) \\
\bar{w}_{3} \leq \sqrt{\bar{Q}}=118.144 & \text { (vs. } \left.7.8544^{2}=61.69\right)
\end{array}
$$

where $w_{i}=\bar{w}_{i} \sqrt{E I}$

HW \#9 Solution
(1) ${ }^{\text {a) }} u^{\prime \prime}+\frac{2}{\rho} u^{\prime}+u=0, \quad 1=\frac{d}{d \rho}$
b) $u=a_{0}+a_{1} \rho+a_{2} \rho^{2}+\cdots$

Substitute, collect tums, set coefficient of $f^{n}=0$
Fid $a_{1}=0$ and all $a_{o d d}=0$

$$
\begin{aligned}
& a_{2}=\frac{-a_{0}}{6}, \quad a_{4}=\frac{a_{0}}{120}\left(=-\frac{a_{2}}{20}\right) \\
& a_{6}=-\frac{a_{4}}{42}=-\frac{a_{0}}{5040}=-\frac{a_{0}}{7!}, a_{8}=\frac{a_{0}}{9!} \\
& u(\rho)=a_{0}\left(1-\frac{\rho^{2}}{3!}+\frac{\rho 4}{5!}-\frac{\rho}{7!}+\frac{\rho^{8}}{9!}+\cdots\right) \\
& \left(=a_{0} \frac{\sin \rho}{\rho}\right)
\end{aligned}
$$

c)

$$
\begin{aligned}
& \frac{d u(\rho)}{d \rho}=0=a_{0}\left(-\frac{2 \rho}{3!}+\frac{4 \rho^{3}}{5!}-\frac{6 \rho^{5}}{7!}+\frac{8 \rho^{\frac{7}{7}}}{9!}+\cdots\right) \\
& \rho=0 \text { and }-\frac{\rho}{3}+\frac{\rho^{3}}{30}-\frac{\rho^{5}}{840}+\frac{\rho^{7}}{45360}-\cdots=0
\end{aligned}
$$

a root solver gives $\rho=4.14$

$$
\Rightarrow \quad \rho=\frac{w_{1} R}{c}=4.4, \quad w_{1}=\frac{4_{1} 14 c}{R}
$$

Lord Raylagh (1872) gives $1.43 \pi=4.49$
(2) $x^{2} J_{0}^{\prime \prime}+x J_{0}^{\prime}+x^{2} J_{0}=0$
diviletg $x^{2}: J_{0}^{\prime \prime}=-\frac{J_{0}^{\prime}}{x}-J_{0}$
$D_{1}$ ffientiate $\Rightarrow J_{0}^{\prime \prime \prime}=-\frac{J_{0}^{\prime \prime}}{x}+\frac{J_{0}^{\prime}}{x^{2}}-J_{0}^{\prime}$
Muct by $x^{2}: \quad x^{2} J_{0}^{\prime \prime \prime}+x J_{0}^{\prime \prime}-\left(1-x^{2}\right) J_{0}^{\prime}=0$
Let $f=-J_{0}^{\prime}$

$$
\left.\begin{array}{rl}
-x^{2} f^{\prime \prime}-x f^{\prime}+\left(1-x^{2}\right) f & =0 \\
x^{2} f^{\prime \prime}+x f^{\prime}+\left(x^{2}-1\right) f & =0
\end{array}\right\}
$$

$$
\left.J_{1} \text { satisfeis } \quad x^{2} J_{1}^{\prime \prime}+x J_{1}^{\prime}+\left(x^{2}-1\right) J_{1}=0\right\}
$$

Compasion of eqs $\Rightarrow J_{1}$ and $f$ satigy the same $O D E$
Both Jo and $J_{1}$ ODE'SA adpit ino ineasly indyemenent Solutions, one bounded as $x \rightarrow 0$, one unbourded.

Sinie both $J_{0}$ \& $J_{1}$ are bounded, we have that $f$ and $J_{1}$ are at least a multigle of ore another. Normalezation givis

$$
\begin{aligned}
& f=J_{1}, B_{u} t f=-J_{0}^{\prime} \\
& \therefore-J_{0}^{\prime}=J_{1}
\end{aligned}
$$

HT \# 10 Solutions

1. $\frac{d^{2} x}{d t^{2}}+x=\alpha x^{5}$
$\operatorname{Set} x=A \cos n t$
$-A \omega^{2} \cos \omega t+A \cos \omega t=\alpha A^{5} \cos ^{5} \omega t$
Identity (maxima) : $\cos ^{5} \theta=\frac{5}{8} \cos \theta+\frac{5}{16} \cos 3 \theta+\frac{1}{16} \cos 5 \theta$
So $\alpha A^{5} \cos ^{5} \omega t=\frac{5}{8} A^{5} \cos \omega t+\underset{\substack{\text { nonresonant } \\ \text { terms }}}{ }$

$$
\begin{aligned}
& -A \omega^{2}+A=\frac{5}{8} A^{5} \alpha \\
& \Rightarrow \quad \omega^{2}=1-\frac{5}{8} \alpha A^{4}
\end{aligned}
$$

2. $\ddot{x}+x=0.1\left(1-2 x^{2}+6 x^{4}\right) \dot{x}$

2a. $x=A \cos \omega t$

$$
-A \omega^{2} \cos \omega t+A \cos \omega t=0.1\left(1-2 A^{2} \cos ^{2} \omega t+b A^{4} \cos ^{4} \omega t\right) * A
$$

$$
-A \omega \sin \omega t
$$

$$
\begin{aligned}
& \text { Identities (maxima): } \cos ^{2} \theta \sin \theta=\frac{1}{4} \sin \theta+\frac{1}{4} \sin 3 \theta \\
& \cos ^{4} \theta \sin \theta=\frac{1}{8} \sin \theta+\frac{3}{16} \sin 3 \theta+\frac{1}{16} \sin 5 \theta \\
&\left(-\omega^{2}+1\right) A \cos \omega t=0.1(-A \omega)(\sin \omega t)\left(1-2 A^{2}\left(\frac{1}{4}\right)+6 A^{4}\left(\frac{1}{8}\right)\right)
\end{aligned}
$$

thonresonant tums

Balancing the harmonies:
$\cos \omega t:\left(-\omega^{2}+1\right) A=0 \Rightarrow \omega=1$
sinnt: $A w\left(1-\frac{A^{2}}{2}+b \frac{A^{4}}{8}\right)=0$

$$
A^{2}=\frac{2}{b}(1 \pm \sqrt{1-2 b})
$$

2b. For $b>\frac{1}{2}$ there me no real coots, no LC's

$$
\begin{array}{r}
b<\frac{1}{2} \quad \text { " real positine poots } \\
(=2 \text { LC's })
\end{array}
$$

7or $b=\frac{1}{2}$ there is one degencrate $\angle C$

2c. Check with "pplane".
For $b=\frac{1}{4}$, theay gives $A^{2}=8 \pm 2^{5 / 2}=2.34,13.65$

$$
\Rightarrow A=1.53,3.70
$$

Agnees w/ pplane plot:

3.

$$
\begin{aligned}
& \dot{x}=y+1 x^{3}+\alpha x \\
& \dot{y}=-x+.1 \dot{x}^{3}-\beta \dot{x}
\end{aligned}
$$

Ba. Differentiate $i^{\text {st }}$ equation:

$$
\begin{aligned}
\ddot{x} & =\dot{y}+(.1)\left(3 x^{2} \dot{x}\right)+\alpha \dot{x} \\
& =-x+(.1) \dot{x}^{3}-\beta \dot{x}+(.1)\left(3 x^{2} \dot{x}\right)+\alpha \dot{x}
\end{aligned}
$$

02

$$
\ddot{x}+x=(\alpha-\beta) \dot{x}+(.1)\left(\dot{x}^{3}+3 x^{2} \dot{x}\right)
$$

3b. Harmonic balance: $x=A \cos \omega t$

$$
\begin{aligned}
-\omega^{2} A \cos \omega t+A \cos \omega t & =(\alpha-\beta)(-A \omega \sin \omega t) \\
& +(1)\left(A^{3}\right)\left(-\omega^{3} \sin ^{-3} \omega t\right. \\
& \left.-3 \cos ^{2} \omega t \sin \omega t\right)
\end{aligned}
$$

$\sin ^{3} \omega t=\frac{3}{4} \sin \omega t-\frac{1}{4} \sin 3 \omega t$
$\cos ^{2} \omega t \sin \omega t=\frac{1}{4} \sin \omega t+\frac{1}{4} \sin 3 \omega t$
$\cos \omega t:\left(-\omega^{2}+1\right) A=0 \Rightarrow \omega=1$
sin $\omega$ t: $O=(\alpha-\beta)(-A \omega)+(.1) A^{3}\left(-\frac{3}{4} \omega^{3}-\frac{3}{4}\right) \Rightarrow$
(Using $\omega=1$ ) $A^{2}=\frac{20}{3}(\beta-\alpha)$
72 a real solution, $A^{2}>0 \Rightarrow \beta>\alpha$
Example $\beta=.2, \alpha=.1, A^{2}=\frac{2}{3}, A \approx .82$
agrees w/ simulation on plane 6

Bc, A Hop bifurcation occurs for $\beta=\alpha$
The LC exists for $\beta>\alpha$.
What is the stability of the origin for $\beta>\alpha$ ? The $O D E$ becomes (linearize hear the ongin);

$$
\ddot{x}+x=(\alpha-\beta) \dot{x}+\text { nonlinear farms }
$$

for $\beta>\alpha$ this is a damping tam
$\Rightarrow$ the origin is stable
$\Rightarrow$ The LC is unstable:


Motions near the LC move away from it and head fowends the origin
$\therefore$ The Hop bifurcation is SUBCRITICAL
(Thesis agrees with simulation using plane, Undu "OPTIONS" Choise"sOLUTION DIRECTION" as forward. You will see the LC as repelling (ie. unstable.)

