HWI SOLUTIONS



Small & approx



$$\frac{d}{dt} \frac{\partial L}{\partial \dot{o}} - \frac{\partial L}{\partial 0} = 0$$

$$\frac{d}{dt} \left(2 \left(\frac{1}{a} (\Xi_{6})_{p}^{+} \frac{q}{2} a^{2} M \right) \dot{O} \right] + 5 k a^{2} O = 0$$

$$\left((\Xi_{6})_{p}^{+} q a^{2} M) \dot{O} + 5 k a^{2} O = 0$$

$$\left(\frac{7}{3} m q^{2} + q M q^{2} \right) \ddot{O} + 5 k q^{2} O = 0$$

$$\left(\frac{7}{3} m q^{2} + q M q^{2} \right) \ddot{O} + 5 k q^{2} O = 0$$

$$\left(\frac{3}{3} m q^{2} + q M q^{2} \right) O = 0$$

Note: For KE of beam, can alter naturely use

(3) Use small & approx. Do not neglect gravity. $I_{beam} = \frac{1}{12}m(3a)^2$ T $M = \frac{1}{2}M(3aO)^2 - Mg3a\frac{O^2}{2}$ Shm Fim beam 1 Ibeam 02 - Mg 3 a 02 Von= 3 a 0 + 1 m V cm² Shatom Springs 0 0 2.1 K(200)2 PE loss by rotation of Shm=(3.a). (1-coso)

beamby AOMass M: Mg. Ohm $\Delta h_m = (\frac{3}{2}a) \cdot (1 - \cos 0)$ beam $M \cdot M \cdot g \cdot \Delta h_m$ $Small A \rightarrow \cos 0 \approx 1 - \frac{0^2}{2}$ $\Delta h_m = 3a\frac{0^2}{2}$ $\Delta h_m = \frac{3}{2}a\frac{0^2}{2}$

$$L = T - V$$

$$= \frac{9}{2}Ma^{2}\dot{\Theta}^{2} + \frac{9}{a_{4}}ma^{2}\dot{\Theta}^{2} + \frac{1}{2}m\left(\frac{9}{4}\right)a^{2}\dot{\Theta}^{2} + Mg_{3}a\frac{\Theta^{2}}{a} + Mg_{3}a\frac{\Theta^$$

$$Lagrange = \frac{\lambda}{\lambda t} \frac{\partial L}{\partial 0} - \frac{\partial L}{\partial 0} = 0$$

$$D = \frac{\lambda}{\lambda t} \left[\left(\frac{q}{ma^{2}} + \frac{3}{4}ma^{2} + \frac{q}{4}ma^{2} \right) \right] - \left(\frac{Mg}{3} \frac{3}{4} + \frac{mg}{3} \frac{3}{4} - \frac{8ka^{2}}{6} \right) 0$$

$$0 = (9Ma + 3ma)\ddot{0} + (-Mg3 - Mg\frac{3}{2} + 8ka)0$$

Effective Stiffness is coefficient in front of O (as in mX+Kx=0) System is unstable (by divergence) if effective stiffness <0

Stability:

8ka - (BMg+ 3mg)>0

8ka > 3Mg+3mg $k > (3Mg + \frac{3}{2}mg)$ 8a

In order to see how the $x = c_1 \cos \omega t + c_2 \sin \omega t$ solution agrees with the x = 1 + t solution in the limit as ω approaches zero, we observe that the initial conditions associated with x = 1 + t are:

$$\begin{cases} x(0) = 1\\ x'(0) = 1 \end{cases}$$

applying these initial conditions to $x = c_1 \cos \omega t + c_2 \sin \omega t$, we get:

$$\begin{cases} x(0) = 1 \\ x'(0) = 1 \end{cases} \Rightarrow \begin{cases} c_1 = 1 \\ \omega c_2 = 1 \Rightarrow c_2 = \frac{1}{\omega} \end{cases}$$

so the solution becomes:

$$x = \cos \omega t + \frac{1}{\omega} \sin \omega t$$

now, we take the limit of this solution as $\omega \to 0$:

$$\lim_{\omega \to 0} x = \lim_{\omega \to 0} \cos \omega t + \lim_{\omega \to 0} \frac{1}{\omega} \sin \omega t = 1 + \lim_{\omega \to 0} \frac{t \cos \omega t}{1} = 1 + t$$

where we use l'hopital's rule to evaluate the second limit.

In summary, since one of the arbitrary constants actually depends on ω , it is essential to account for it when taking the limit $\omega \to 0$.



So Lagrange's equation is

$$f'' + f''_{\frac{x^2}{2}} + kx = 0 \qquad Ans: to part a$$

where
$$f(x) = m_1 + \frac{(d-x)^2 m_2}{a^2 - (d-x)^2}$$

b) To find natural frequency, keep only linear terms. This gives $f(0)\dot{x} + kx = 0$

$$\omega^{2} = \frac{k}{f(o)} = \frac{k}{\frac{d^{2}m_{z}}{a^{2}-d^{2}}} + m_{1}$$



#2 continued
Similar to problem #1,

$$\frac{d}{dt} \frac{2t}{2x} - \frac{2t}{3x} = 0 \implies f \stackrel{"}{x} + f \stackrel{'}{x} \stackrel{"}{z} + 2mg = 0$$
where $f(x) = m(1+4x^2)$
lirearize around $x=p$:
 $f(0)\stackrel{"}{x} + 2mg = 0$, $f(0) = m$
 $\stackrel{"}{x} + 2g = 0$
 $w^2 = 2g$, $w = \sqrt{2g}$
#3. $x_p = R \cos(52t-\phi)$
 $x_p = -R \Omega \sin(52t-\phi)$
 $amglitude ef veloity = R S\Omega$
thom class notes, $\frac{R}{5st} = \frac{1}{\sqrt{(1+\frac{1}{3t})^2 + 4\frac{1}{3t}}}$
 $\frac{R \Omega}{5st} = \frac{\Omega}{\sqrt{(1+\frac{1}{3t})^2 + 4\frac{1}{3t}}}$
 $\frac{R \Omega}{\sqrt{(1-\frac{1}{3t})^2 + 4\frac{1}{3t}}} = \frac{2}{\sqrt{(1-\frac{1}{3t})^2 + 4\frac{1}{3t}}}$

Set
$$\frac{d}{dg}\left(\frac{g}{\sqrt{(1-q^2)^2+4p^2q^2}}\right) = 0$$
 for max

Obtain
$$\frac{q^{q}-1}{((1-q^{q})^{2}+4\beta^{2}q^{2})^{3/2}} = 0 \implies q=1$$

... Max velocity amplitude occurs at
$$[S2=w]$$

#4.
$$\theta = \sqrt{\frac{2mgR}{mR^2 + I_0}} \sqrt{\theta}, \theta(0) = 0$$

has two exact solutions:
One of them is
$$\theta \equiv 0$$
 as stated.
The other is obtained by separation of variables

$$\frac{d\theta}{V\theta} = \sqrt{\cdots} dt$$

$$2\theta'^{2} = \sqrt{\cdots} t + Q^{T} (C=0 \text{ for } \theta(0)=0)$$

$$\therefore \theta(t) = \frac{1}{2}\sqrt{\cdots} t^{2}$$

#4 continued

Comment: An equation with two (or more) solutions is bad news: how do you know if the solution you have found is the right one for your application?

That is why the mathematicians have given us a uniqueness proof which gives conditions on the right hand side of The differential equation (f(0)) $\frac{d\theta}{dt} = f(\theta)$

Which guarantee that there is only one solution.



(b) Linearise about
$$x=0=0$$
 ($\cos 0 \ge 1 \sin 0 \ge 0$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$
 $x) [max + ma^{2} \frac{1}{2} + mga \frac{0}{2} = 0]$
(max + ma^{2} \frac{1}{2} + mga \frac{0}{2} = 0]
(max + ma^{2} \frac{1}{2} + mga \frac{0}{2} = 0]
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(max + ma^{2}) = 0 = 2 = [\frac{1}{0}] = 2 = [\frac{1}{0}] = 2 = [\frac{1}{0}] = 0
 $\frac{1}{2} = [\frac{1}{0} + mga \frac{1}{2}] = 0$
(1-2 ω^{2})(1- ω^{2}) = $\omega^{4} = 0$
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$$\begin{array}{c} (\bigcirc X_{3}^{T}MX_{1}=0 \\ X_{3}^{T}MX_{1}=0 \\ X_{3}^{T}KX_{1}=0 \\ X_{3}^{T}KX_{1}=0 \\ X_{3}^{T}KX_{1}=0 \\ Soluting in MATLAB, all conditations of above equations yield zero. \\ (\textcircled{black}) R = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \\ R^{T}MR = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \\ diagonal \\ (\textcircled{black}) X_{3} \otimes as | inter combinations of privilations coordinates p_{1} > Pa \\ principal coordinates p = P^{-1}Z \\ (Pg 181) \\ model matrix P = \begin{bmatrix} 1 & 2 \\ 0 & 1 \end{bmatrix} \\ X = -0.5257 p_{1} - 0.8507 p_{2} \\ (\varTheta{black}) = -0.3249 p_{1} + 1.3764 p_{2} \\ (\rule{black}) R^{T}MR \stackrel{\circ}{p} + R^{T}KR p = 0 \\ \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \stackrel{\circ}{p} + \begin{bmatrix} 0.382 & 0 \\ 0 & 2.6180 \end{bmatrix} p = 0 \\ (\rule{black}{} 0 & 2.6180 \end{bmatrix} p = 0 \\ \end{array}$$

(1)
$$M\ddot{x} + 3Kx + K(\frac{1}{a}x - \sqrt{\frac{3}{2}}y) + \frac{1}{a}K(\frac{1}{a}x + \sqrt{\frac{3}{2}}y) = 0$$

 $M\ddot{x} + X(3K + \frac{1}{a}K + \frac{1}{4}K) + y(-\frac{13}{a}K + \frac{\sqrt{3}}{a}, \frac{1}{a}K)$
 $\rightarrow M\ddot{x} + \frac{\sqrt{5}}{4}Kx - \sqrt{\frac{3}{4}}Ky = 0$
(2) $M\ddot{y} + K(-\frac{13}{2}x + \frac{3}{a}y) + \frac{1}{a}K(\frac{\sqrt{3}}{a}x + \frac{3}{a}y) = 0$
 $\rightarrow M\ddot{y} - \frac{13}{4}Kx + \frac{9}{4}Ky = 0$
 $M = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \quad K = \begin{bmatrix} \frac{15}{4} & -\frac{\sqrt{3}}{4}y \\ -\frac{\sqrt{3}}{4}y & \frac{9}{4}y \end{bmatrix}$ (setting $k = 1$)
(setting $m = 1$)

$$MATLAB [v,d] = eig(H,M) d = [w_1^2 0] w_1 = 1.4608 w_a = 1.9662$$



Expand V using the Taylor series for
$$\sqrt{1+Z}$$

 $\sqrt{1+Z} = 1+\frac{Z}{2} - \frac{Z^2}{8} + \cdots$
 $\sqrt{(\frac{1}{2}-x)^2 + (\frac{\sqrt{3}}{2}-y)^2} = \sqrt{\frac{1}{4}-x+x^2+\frac{3}{4}-\sqrt{3}y+y^2}$
 $= \sqrt{1+(-x+x^2-\sqrt{3}y+y^2)}$
 $= 1+(\frac{-x+x^2-\sqrt{3}y+y^2}{2})$
 $= (-\frac{x+x^2-\sqrt{3}y+y^2}{2})^2 + \cdots$
 $= 1-\frac{y}{2}-\sqrt{\frac{3}{2}y}+\frac{x^2}{2}+\frac{y^2}{2}$
 $-\frac{1}{8}(x^2+3y^2+2\sqrt{3}xy)+x$
 $= 1-\frac{y}{2}-\frac{\sqrt{3}}{2}y+\frac{3}{8}x^2+\frac{1}{8}y^2-\frac{\sqrt{3}}{4}xy$
 $(\sqrt{(\frac{1}{2}-x)^2+(\frac{\sqrt{3}}{2}-y)^2-1)^2} = (\frac{1}{2}-x)^2+(\frac{\sqrt{2}}{2}-y)^2+1-2(\sqrt{4}))$
 $= \frac{1}{4}-x+x^2+\frac{3}{4}-\sqrt{3}}y+y^2+1$
 $-2+x+\sqrt{3}}y-\frac{3}{4}x^2-\frac{1}{4}y^2+\frac{\sqrt{3}}{2}xy$
 $= x^2+y^2-\frac{3}{4}x^2-\frac{1}{4}y^2+\frac{\sqrt{3}}{2}xy$
 $= \frac{1}{4}x^2+\frac{3}{4}y^2+\frac{\sqrt{3}}{2}xy$

$$\sqrt{(1+x)^{2} + y^{2}} = \sqrt{1+2x+x^{2}+y^{2}}$$

$$= 1 + \frac{1}{2} (2x+x^{2}+y^{2}) - \frac{1}{8} (2x+x^{2}+y^{2})^{2} + x^{2}$$

$$= 1 + x + \frac{x^{2}}{2} + \frac{y^{2}}{2} - \frac{x^{2}}{2}$$

$$= 1 + x + \frac{y^{2}}{2} - \frac{x^{2}}{2}$$

$$= 1 + x + \frac{y^{2}}{2} - \frac{x^{2}}{2}$$

$$= 1 + x + \frac{y^{2}}{2} - \frac{x^{2}}{2}$$

$$= 1 + 2x + x^{2} + y^{2} + 1 - 2(1 + x + \frac{y^{2}}{2})$$

$$= 1 + 2x + x^{2} + y^{2} + 1 - 2 - 2x - y^{2}$$

$$= x^{2}$$

$$V = \frac{1}{2}k\left(\frac{x^{2}}{4} + \frac{3}{4}y^{2} + \sqrt{\frac{3}{2}}xy\right)$$

+ $\frac{1}{2}(2k)\left(\frac{x^{2}}{4} + \frac{3}{4}y^{2} - \sqrt{\frac{3}{2}}xy\right)$
+ $\frac{1}{2}(3k)x^{2}$
$$V = \frac{1}{2}k\left[\chi^{2}\left(\frac{1}{4} + \frac{2}{4} + 3\right) + \gamma^{2}\left(\frac{3}{4} + \frac{2}{4}\right) + \sqrt{3}xy\left(\frac{1}{2} - 1\right)\right]$$

= $\frac{1}{2}k\left[\frac{1/\frac{4}{7}}{7}x^{2} + \frac{9}{4}y^{2} - \sqrt{\frac{3}{2}}xy\right]$
 $m_{X}^{2} = -\frac{2V}{2x} = -\frac{15}{4}kx + \sqrt{\frac{3}{7}}yk$
 $m_{Y}^{2} = -\frac{2V}{2y} = -\frac{9}{4}ky + \frac{\sqrt{3}}{7}xk$

$$M \ddot{x} + K x = 0$$

$$\begin{bmatrix} m & 0 \\ 0 & m \end{bmatrix} \begin{bmatrix} \ddot{x} \\ \ddot{y} \end{bmatrix} + \begin{bmatrix} \frac{12}{5}k & -\frac{\sqrt{3}}{4}k \\ -\sqrt{3}k & \frac{9}{4}k \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$x = A \cos \omega t$$

$$y = B \cos \omega t$$

$$\begin{vmatrix} -\omega^{2}m + \frac{15}{4}k & -\frac{\sqrt{3}}{4}k \\ -\frac{\sqrt{3}}{4}k & -\frac{\sqrt{3}}{4}k \end{vmatrix} = 0$$

$$\left(-\frac{\beta^{2}}{7} + \frac{15}{7} \right) \left(-\frac{\beta^{2}}{7} + \frac{9}{7} \right) - \frac{3}{16} = 0 \quad \text{where } \beta = \omega \sqrt{\frac{m}{k}}$$

$$\beta^{4} - 6\beta^{2} + \frac{33}{4} = 0$$

$$\beta^{2} = \frac{6 \pm \sqrt{36-33}}{2} = 3\pm \frac{\sqrt{3}}{2}$$

$$\omega = \sqrt{3}\pm \frac{\sqrt{3}}{2} \sqrt{\frac{k}{m}} = 1.461 \sqrt{\frac{k}{m}}, 1.966 \sqrt{\frac{k}{m}}$$

3. The general motion of the first coordinate of a two degree of freedom system is given by:

$$x_1(t) = R_1 \cos(\omega_1 t - \theta_1) + R_2 \cos(\omega_2 t - \theta_2)$$

Is this a periodic motion? Under what condition will it be periodic?

At t = 0,

$$x_1(0) = R_1 \cos(\theta_1) + R_2 \cos(\theta_2)$$

At what time t will this happen again? Suppose that $\omega_2 = \frac{m}{n}\omega_1$, where m and n are whole numbers. Then

$$x_1(t) = R_1 \cos(\omega_1 t - \theta_1) + R_2 \cos(\frac{m}{n}\omega_1 t - \theta_2)$$

After time $T = \frac{2\pi n}{\omega_1}$, we have

$$x_1(T) = R_1 \cos(2\pi n - \theta_1) + R_2 \cos(2\pi m - \theta_2) = x_1(0)$$

In fact, $x_1(T+t) = x_1(t)$ for all t, not just t = 0. Thus in this case the motion is periodic.

However, if the ratio of ω_2 to ω_1 is an irrational number, then $x_1(t)$ will never return to $x_1(0)$ and the motion will not be periodic.

$$\begin{aligned} \int ct \ X = Rp, \quad R = \begin{bmatrix} I \ I, 2 \ X \end{bmatrix} = \begin{bmatrix} I & I \\ I & -I \end{bmatrix} \\ X_{1} = p_{1} + p_{2} \\ X_{2} = p_{1} - p_{2} \end{aligned}$$

$$\begin{aligned} R^{t} MR \ p^{r} + R^{t} KR p = R^{t} f(t) = \begin{bmatrix} I & I \\ I - I \end{bmatrix} \begin{bmatrix} 0 \\ F \end{bmatrix} \text{ on } \Omega t \\ \begin{bmatrix} I & I \\ I - I \end{bmatrix} \begin{bmatrix} I & I \\ I - I \end{bmatrix} + \begin{bmatrix} I & I \\ I - I \end{bmatrix} \begin{bmatrix} 2 & -I \\ -I & 2 \end{bmatrix} \begin{bmatrix} I & I \\ I - I \end{bmatrix} = \begin{bmatrix} 2 & 0 \\ 0 & c \end{bmatrix} \\ \begin{bmatrix} I & 3 \\ I & -3 \end{bmatrix} \\ p_{1} + W_{1}^{2} p_{1} = \sum co_{S} \Omega t, \quad W_{1} = I \\ p_{2}^{2} + W_{2}^{2} p_{2}^{2} - \sum co_{S} \Omega t, \quad W_{2} = V3 \end{aligned}$$

$$p_{1} = K \cos \Omega t, \quad (-\Omega^{2} + W_{1}^{2})K = F \\ So \qquad p_{1} = \frac{F/2}{I - \Omega^{2}} \cos \Omega t \\ Similarly \qquad p_{2} = \frac{-F/2}{3 - \Omega^{2}} \cos \Omega t \\ \therefore \quad X_{1} = p_{1} + p_{2} = \left(\frac{I}{I - \Omega^{2}} - \frac{1}{3 - \Omega^{2}} \right) \sum co_{S} \Omega t \\ X_{2} = p_{1} - p_{2} = \left(-\frac{I}{I - \Omega^{2}} + \frac{1}{3 - \Omega^{2}} \right) \sum co_{S} \Omega t \\ = \frac{(2 - \Omega^{2})F}{(I - \Omega^{2})(3 - \Omega^{2})} \cos \Omega t \end{aligned}$$

$$\dot{X}_{1} + 2X_{1} - X_{2} = 0$$

$$\dot{X}_{2} - X_{1} + 2X_{2} = F \cos \Omega t$$
Sut $X_{1} = A \cos \Omega t$

$$X_{2} = B \cos \Omega t$$

$$-\Omega^{2}A + 2A - B = 0$$

$$-\Omega^{2}B - A + 2B = F$$
Solve for A, B

$$\Rightarrow A = \frac{F}{(1 - \Omega^{2})(3 - \Omega^{2})}, \quad B = \frac{F(2 - \Omega^{2})}{(1 - \Omega^{2})(3 - \Omega^{2})}$$
Agrees with (1)

SOLUTION to question 3:

Multiply the first eq. in (7) by $-\Omega^2$ and add to the second eq. in (7) giving:

$$R^t(-\Omega^2 M + K)R = -\Omega^2 D_1 + D_2$$

Take the inverse of both sides:

$$(R^{t}(-\Omega^{2}M+K)R)^{-1} = (-\Omega^{2}D_{1} + D_{2})^{-1}$$
$$R^{-1}(-\Omega^{2}M+K)^{-1}(R^{t})^{-1} = (-\Omega^{2}D_{1} + D_{2})^{-1}$$

Now multiply on the left by R and on the right by R^t , giving

$$(-\Omega^2 M + K)^{-1} = R(-\Omega^2 D_1 + D_2)^{-1} R^t$$

This demonstrates the equivalence of eqs.(5) and (13).



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$$C) \quad \mathcal{N}(\mathbf{x}, \mathbf{t}) = \begin{cases} \sum_{n=1}^{\infty} (A_n \cos \omega_n t + b_n \sin \omega_n t) \cos \omega_n \mathbf{x} \\ + A_0 + b_0 t \\ & \int \mathcal{N}_{ij} d body \ mode \end{cases}$$

$$d) \quad \mathbf{IC} \quad t = 0, \ \mathcal{U}_t = 0 \quad \Rightarrow \quad b_n = 0, \ n = 0, \ l, 2, \dots \\ t = 0, \ \mathcal{U}_t = 0 \quad \Rightarrow \quad b_n = 0, \ n = 0, \ l, 2, \dots \\ t = 0, \ \mathcal{U}_t = A_0 + \quad \sum_{n=1}^{\infty} A_n \cos \frac{n\pi t}{t} \\ Mull by \quad \cos \frac{m\pi x}{t} \quad d \int_0^d \Rightarrow \\ & \int_0^d \frac{1}{t} \cos \frac{m\pi x}{t} dx = \int_0^d A_n \left(\cot \frac{m\pi x}{t} \right)^2 d\mathbf{x} = \frac{d}{t} a_n \\ (m > 0) \\ A_m = \frac{2}{t} \int_0^d \frac{1}{t} \cos \frac{m\pi x}{t} dx, \ m > 0 \\ A_0 = \frac{1}{t} \int_0^d \frac{1}{t} d\mathbf{x} = \frac{1}{t^2} \left[\frac{x^2}{t} \Big|_0^d \right] = \frac{1}{2} \\ A_{m} = \begin{cases} \frac{1}{t} & m = 0 \\ 0 & m = 1, \ y \in y_{1,1,1}^{\infty}, \\ -\frac{4}{m^2 \pi^2}, \ m = 1, \ y \in y_{1,2}^{\infty}, \\ -\frac{4}{t} & m^{m} \pi^2, \end{cases}$$

$$e) \quad Cat \quad x = \frac{d}{t}, \ u(d_{t}, t) = \frac{1}{t} + \sum_{n=1,n,r}^{\infty} a_n \cos \frac{n\pi ct}{t} \\ At \quad x = l, \ u(d_{t}, t) = \frac{1}{t} + \sum_{n=1,n,r}^{\infty} a_n \cos \frac{n\pi ct}{t} \\ a_t \quad x = l, \ u(d_{t}, t) = \frac{1}{t} + \sum_{n=1,n,r}^{\infty} a_n \cos \frac{n\pi ct}{t} \\ = \frac{1}{t} - \sum_{n=1,n,r}^{\infty} a_n \cos \frac{n\pi ct}{t} \end{cases}$$

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$$f) \quad n(o,t) = \pm - \pm - \pm \left(\cos \frac{\pi}{2} t + \frac{1}{9} \cos \frac{\pi}{2} t + \frac{1}{7} \cos \frac{\pi}{2} t + \cdots \right)$$

See figure attached.
2.
$$4$$
 $x=0$ $x=1$
 $EA \frac{\partial u}{\partial x} = -ku$ at $x=1$
 $Aee text, problem 7.8, pp. 210-211$
3. 4 M
 $x=0$ $x=1$
 $EA \frac{\partial u}{\partial x} = -M \frac{\partial^2 u}{\partial t^2}$ at $x=1$
 $pee text, problem 7.4, p. 207$



Here
$$A_n = -\frac{F_n}{(n\pi)^{-1}}$$

 $B_n = \frac{H_n}{n\pi}$
where F_n and H_n are given in b) and c)
f) $M(x_1t) \approx p_1(t) U_1^-(x) + p_0(t) U_0(x)$
 $\approx (A_1 \cos \pi t + B_1 \sin \pi t + \frac{F_1}{T_1} \cos t) \cos \pi x + p_0(t)$
 $\approx (-\frac{4}{\pi^n(\pi^{-1})} (\cos t - \cos \pi t) - \frac{4}{\pi^3} \sin \pi t) \cos \pi x + p_0(t)$
 $\approx (-\frac{4}{\pi^n(\pi^{-1})} (\cos t - \cos \pi t) - \frac{4}{\pi^3} \sin \pi t) \cos \pi x + \frac{1}{2}t - \frac{1}{3}\cos t + \frac{1}{3}$
2a) $M = U(x) \cos \omega t$
 $-\omega^{\alpha} U + \frac{E_1}{S} U^{-1V} = 0$
 $U^{1V} - k^{\alpha} U = 0, \quad k^{\alpha} = \omega^{2} \frac{g}{E_1}$
 $U = e^{Ax} \Rightarrow \lambda^{\alpha} - k^{\alpha} = \Rightarrow \lambda = k_1 - k_1 \cdot k_2 \cdot k_1$
 $U = c_1 \cosh kx + c_2 \sinh kx + c_3 \cosh kx + c_4 \sinh kx$
BC $U(\alpha) = U(2) = U'(\alpha) = U'(2) = 0$
 q homogeneous algebraic eqs.
For nontrivial solution, set determinant = 0
while gives
 $\cos kl \cosh kl = l$, $k = \sqrt{\omega} (\frac{g}{E_1})^{1/4}$
 $\omega = (kl)^{2} (\frac{E_1}{S})^{1/2}$

Solving the system of 4 homog. eqs., obtain

$$U_n(x) = \cosh(kd) \frac{x}{L} - \cosh(kd) \frac{x}{L}$$

 $+ \mu (\sinh(kd) \frac{x}{L} - \sinh(kd) \frac{x}{L})$

where
$$\mu = -\frac{(\cosh kl - \cos kl)}{(\sinh kl - \sin kl)}$$

b) kl = 4,73, 7.85, 10.99, 14.13

At the bottom of the Table we find: $\omega = \frac{\lambda^2}{l^2} \sqrt{\frac{EI}{S}} \quad \text{where } \lambda = k\ell \text{ in above notation} \\ \frac{\lambda^2}{l^2} \sqrt{\frac{EI}{S}} \quad \text{and } g \text{ is listed as puin Table.}$ $J(u) = \cosh u - \cos u \\ H(u) = \sinh u - \sin u$ Using this notation, the Table gives the mode shape as $J(u) = \int \frac{1}{2} \int \frac{1}{2$

$$U_n(x) = J(\lambda_n \frac{x}{2}) - \frac{J(\lambda_n)}{H(\lambda_n)} H(\lambda_n \frac{x}{2})$$



 $\frac{d^2u}{dx^2} + u = 1$ 3. u(o)=0 2(1)=0 General solution U=Asmx+Bcosx+1 u(o) = B+1 =0 ⇒ B=-1 U(IT) = -BH = 0 => B=1 Since B cannot be equal to both 1 and -1, this problem has no solution. Note that the above system does have a solution for appropriate choices of the jught hand side. For example $\frac{d^2 u}{dx^2} + u = \cos 3x$ & particular solution $n = Asin \times + Bas \times -\frac{1}{2}a + 3 \times$ $u(\pi) = B - \frac{1}{8} = 0 \implies B = \frac{1}{8}$ $u(\pi) = -B - \frac{1}{8}(-1) = 0 \implies B = \frac{1}{8}$ This can be applained in terms of the "Fredholm alternative therem" which I will go over in class.

Ia.



From Table,
$$\lambda_1 = 3.9266$$

 $\omega_1 = \lambda_1^2 \sqrt{\frac{EI}{SL^4}} = 15.42 \sqrt{\frac{EI}{SL^4}}$

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1b.
$$\frac{d^{4}u}{dx^{4}} = 1$$

 $u = c_{1} + c_{2}x + c_{3}x^{2} + c_{4}x^{3} + \frac{x^{4}}{24}$

BC:
$$X = 0, u = 0, u' = 0 \implies c_1 = c_2 = 0$$

 $X = l, u = 0, u'' = 0 \implies c_3 = \frac{1}{16}, c_4 = -\frac{5}{48} l$

$$u = \frac{x^{2}l^{2}}{16} - \frac{5}{48} \frac{x^{3}l}{24} + \frac{x^{4}}{24} \left(= V(x) \text{ below}\right)$$

$$IC.$$

$$Q = EI \int_{0}^{L} (V'') dx = \frac{EI \frac{L^{5}}{320}}{\int \frac{19 L^{9}}{1451520}} = \frac{EI \frac{L^{5}}{320}}{\int \frac{19 L^{9}}{1451520}}$$

$$w_{1} < \sqrt{Q} = \sqrt{\frac{4536}{19}} \sqrt{\frac{EI}{gL^{4}}} = 15.45 \sqrt{\frac{EI}{gL^{4}}}$$
Constant with 10^{1}

 (\tilde{l})

1d. Again take
$$V(x) = \frac{x^2 \ell^2}{16} - \frac{5}{48} \frac{x^3 \ell}{24} + \frac{x^4}{24}$$



E

$$T(x) = x+1 \qquad (See 7.22 m p 225)$$

$$A(x) = TT(x+1)^{2}$$

$$Q = \frac{\int_{0}^{2} EA(x)(u')^{2} dx}{\int_{0}^{2} gA(x) u^{2} dx}$$

Choose $\mathcal{U}(X) = X(X-2)$ Which satisfies the BC $\mathcal{U}(0) = \mathcal{U}(2) = 0$ Then $\mathcal{U}' = 2X - 2$

$$Q = E \int_{0}^{2} \pi (x+1)^{2} (2x-2)^{2} dx = \frac{184}{15} E$$

$$g \int_{0}^{2} \pi (x+1)^{2} x^{2} (x-2)^{2} dx = \frac{464}{105} S$$

$$Q = \frac{161}{58} = \Rightarrow w_1 \le \sqrt{Q} = 1.67 = \frac{1}{58}$$

2.

$$u'' - u = 1$$

3.

a)
$$\mathcal{U} = c_1 \sin x + c_2 \cos x + c_3 \sinh x + c_4 \cosh x - 1$$

 $\mathcal{U}(0) = 0 = c_2 + c_4 - 1$
 $\mathcal{U}'(0) = 0 = -c_2 + c_4$

$$\begin{cases} \mathcal{U}(\pi) = 0 = c_4 \tilde{c} + c_3 \tilde{s} - c_2 - 1 \quad \text{where } \tilde{c} = \cosh(\pi) \\ \tilde{s} = \sinh(\pi) \end{cases}$$

$$\begin{cases} \mathcal{U}'(\pi) = 0 = c_4 \tilde{c} + c_3 \tilde{s} + c_2 - 1 \quad \text{where } \tilde{c} = \cosh(\pi) \\ \tilde{s} = \sinh(\pi) \end{cases}$$

$$\begin{cases} \mathcal{U}''(\pi) = 0 = c_4 \tilde{c} + c_3 \tilde{s} + c_2 - 1 \quad \text{where } \tilde{c} = \cosh(\pi) \\ \tilde{s} = \sinh(\pi) \end{cases}$$

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$$\begin{cases} \mathcal{U}''(\pi) = 0 = c_4 \tilde{c} + c_3 \tilde{s} + c_2 - 1 \quad \text{where } \tilde{c} = \cosh(\pi) \\ \tilde{s} = \sinh(\pi) \end{cases}$$

Consider a clamped-free beam of constant depth, and a width which varies linearly from a maximum at the fixed end to zero at the free end. Taking the origin of coordinates at the fixed end, $I = I_0(1 - x/l)$ and $\mu = \mu_0(1 - x/l)$, where I_0 and μ_0 are respectively the moment of inertia and mass per unit length at the fixed end. Choose a two-term series,

$$W = A_1 x^2 + A_2 x^3$$

$$W = A_{1}x^{2} + A_{2}x^{3}$$

The kinetic and potential energies are:
$$T^{*} = \frac{1}{2} \int_{0}^{l} \mu W^{2} dx = \frac{\mu_{0}}{2} \int_{0}^{l} \left(1 - \frac{x}{l}\right) (A_{1}x^{2} + A_{2}x^{3})^{2} dx$$
$$V_{\max} = \frac{1}{2} \int_{0}^{l} EI(W'')^{2} dx = \frac{EI_{0}}{2} \int_{0}^{l} \left(1 - \frac{x}{l}\right) (2A_{1} + 6A_{2}x)^{2} dx$$

Integrating, taking the partial derivatives, and substituting into (61.127) gives the pair of equations,

$$(2 - \beta/30)A_1 + (2 - \beta/42)A_2l = 0$$

(2 - \beta/42)A_1 + (3 - \beta/56)A_2l = 0
(61.128a,b)

where $\beta = \mu_0 l^4 \lambda / EI_0$. The roots of the frequency equation are $\beta_1 = 51.25$ and $\beta_2 = 1377$, and, hence, the first two frequencies are $\omega_1^2 \leq 51.25 E I_0 / \mu_0 l^4$ and $\omega_2^2 \leq$ 1377EI 0/40l4

2. Rite on
$$\int_{V=C_1}^{V} x^2 + c_2 \chi^3 + c_3 \chi^4$$

$$\overline{Q} = Q \frac{\rho \ell^4}{EI}, \quad \overline{c_2} = c_2 \ell_1, \quad \overline{c_3} = c_3 \ell^2$$

$$\overline{Q} = \chi^4 \frac{\int_0^{\ell} (\chi^{\prime\prime})^2 dx}{\int_0^{\ell} V^2 dx}$$

$$\int_0^{\ell} (y^{\prime\prime})^2 dx = \int_{\ell} \left(\frac{144}{5}\overline{c_3}^2 + 36\overline{c_2}\overline{c_3} + 16c_1\overline{c_3} + 12\overline{c_2}^2 + 112c_1\overline{c_2} + 4c_1^2\right) = \lambda F$$

$$\int_0^{\ell} V^2 dx = \int_{0}^{\ell} \left(\frac{\overline{c_3}}{4} + \frac{\overline{c_2}\overline{c_3}}{4} + \frac{2}{7}c_1\overline{c_3} + \frac{\overline{c_3}}{7} + \frac{c_1\overline{c_2}}{3} + \frac{c_1^2}{5}\right) = \int_{\ell}^{\ell} G$$

$$\overline{Q} = \frac{F}{G}, \quad G = F$$

$$\frac{1}{44kc} \frac{2}{2i}, \quad \frac{2}{3c_3} + \frac{2}{3c_3} + \frac{1}{3kc_3} +$$

 $\begin{bmatrix} 2\overline{Q} - 40 & \overline{Q} - 36 & 2\overline{Q} - 112 \\ \overline{5} & \overline{3} & \overline{7} & \overline{7} \\ \overline{Q} - 36 & 2\overline{Q} - 168 & \overline{Q} - 144 \\ \overline{7} & \overline{7} & 4 \\ 2\overline{Q} - 112 & \overline{Q} - 144 & 10\overline{Q} - 2592 \\ \overline{7} & 4 & +5 \\ \end{bmatrix} \begin{bmatrix} \zeta_1 \\ \zeta_2 \\ \zeta_3 \end{bmatrix}$

Set $dut=0 \implies$ $5\overline{Q}^{3} - 72324 \overline{Q}^{2} + 35392896 \overline{Q} - 426746880 = 0$ $\overline{Q} = 12,369, +944.322, 13958,107$ Exer: $\overline{w}_{1} \leq \sqrt{\overline{Q}} = 3,517$ (VS. 18\$\$51² = 3.516) $\overline{w}_{2} \leq \sqrt{\overline{Q}} = 22,233$ (VS. 9.694)² = 22,03) $\overline{w}_{3} \leq \sqrt{\overline{Q}} = 118.144$ (VS. 7.8546² = 61.69) where $\overline{w}_{i} = \overline{w}_{i} \sqrt{\frac{FI}{PL^{4}}}$

HW #9 Solution

(i)
$$u'' + \frac{2}{5}u' + u = 0$$
, $i = d$

b)
$$\mathcal{U} = a_0 + a_1 g + a_2 g^* + ...$$

Substitute, collect turns, set coefficient of
$$g^{n=0}$$

Find $a_1 = 0$ and all $a_{odd} = 0$
 $a_2 = -\frac{a_0}{5}$, $a_4 = \frac{a_0}{120}$ $\left(= -\frac{a_2}{20}\right)$
 $a_4 = -\frac{a_0}{42} = -\frac{a_0}{50+0} = -\frac{a_0}{7!}$, $a_8 = \frac{a_2}{9!}$
 $ulgl = a_0 \left(1 - \frac{g^2}{3!} + \frac{g \cdot 4}{5!} - \frac{g^6}{7!} + \frac{g^8}{9!} + \cdots\right)$
 $\left(= a_0 \frac{smg}{s}\right)$
c) $\frac{dugl}{dg} = 0 = a_0 \left(-\frac{2g}{3!} + \frac{4g^5}{5!} - \frac{6g^5}{7!} + \frac{gg^7}{9!} - \frac{g!}{9!}\right)$

$$S^{=0} \text{ and } -\frac{g}{3} + \frac{g^{3}}{30} - \frac{g^{5}}{840} + \frac{g^{7}}{45340} - \cdots = 0$$

a proof solver gives $g = 4.14$
 $\Rightarrow g = \frac{W_{1}R}{c} = 4.14$, $W_{1} = \frac{4.14}{R}$
Lord Rayleigh (1872) gives $1.43\pi = 4.49$

(2)
$$\chi^{2}J_{0}^{"'} + x J_{0}^{'} + \chi^{2}J_{0} = 0$$

divideby χ^{2} : $J_{0}^{"} = -\frac{J_{0}^{"}}{x} - J_{0}^{'}$
Differentiate $\Rightarrow J_{0}^{"'} = -\frac{J_{0}^{"}}{x^{2}} - J_{0}^{'}$
Must by χ^{2} : $\chi^{2}J_{0}^{"'} + xJ_{0}^{"} - (1-\chi^{2})J_{0}^{'} = 0$
Let $f^{2} = -J_{0}^{'}$
 $-\chi^{2}f'' - \chi f^{'} + (1-\chi^{2})f = 0$
 $\chi^{2}f'' + \chi f^{'} + (\chi^{-1})f = 0$
 J_{1} satisfies $\chi^{2}J_{1}^{"} + \chi J_{1}^{'} + (\chi^{-1})J_{1} = 0$
Companion of $ogp \Rightarrow J_{1}$ and f satisfy
the same ODE
Both. J_{0} and J_{1} ODE's two liningly independent
Solutions, one bounded as $\chi \rightarrow 0$, one
unbounded.
Since both. $J_{0} \notin J_{1}$ are at least a multiple
of one construe. Normalization gives
 $f = J_{1}$. But $f = -J_{0}^{'}$
 $\therefore -J_{0}^{'} = J_{1}$

HW # 10 Solutions

1.
$$\frac{d^{2}x}{dt^{2}} + X = d \times 5$$

Set $X = A$ cos wit

$$-Aw^{2}cos wt + A cos wt = a A^{5}cos^{5}wt$$

$$Identity (maxima): cos^{5}\theta = \frac{6}{8}cos \theta + \frac{1}{16}cos 5\theta$$

$$So = a A^{5}cos^{5}wt = \frac{5}{8}A^{5}cos wt + nonresonant + terms$$

$$-Aw^{2} + A = \frac{5}{8}A^{5} \propto$$

$$= W^{2} = 1 - \frac{5}{8} \propto A^{4}$$

2. $\frac{1}{2} + \chi = 0.1(1-2\chi^{2}+b\chi^{4})\dot{\chi}$
2a. $\chi = A cos wt$

$$-Aw^{2}coswt + Acoswt = 0.1(1-2A^{2}cos^{2}wt + bA^{4}cos^{4}wt) * A + \frac{1}{16}cos^{2}wt + \frac{1}{16}cos^{4}wt) * A + \frac{1}{16}cos^{4}wt + \frac{1$$

Balancing the harmonics:

$$(05 \text{ wt}; (-w^{2}+1)A = 0 \implies w = 1$$

 $smwt: Aw(1 - \frac{A^{2}}{2} + b\frac{A^{4}}{8}) = 0 \implies$
 $A^{2} = \frac{2}{b}(1 \pm \sqrt{1-2b})$

2b. For
$$b7\frac{1}{2}$$
 there are no real roots, no LC's
 $b<\frac{1}{2}$ " " 2 real positive roots
 $(= 2 LC's)$
 $for b=\frac{1}{2}$ there is one degenerate LC

2c. Check with "pplane".

$$fr_2 b = \frac{1}{4}$$
, theory gives $A^2 = 8 \pm 2^{5/2} = 2.34$, 13.65
 $\implies A = 1.53, 3.70$
Agrees w/ pplane plot:



3.
$$\dot{x} = y + .1x^{3} + dx$$

 $\dot{y} = -x + .1\dot{x}^{3} - \beta\dot{x}$
3a. Differentiate 1st equation:
 $\ddot{x} = \dot{y} + (.1)(3x^{2}\dot{x}) + d\dot{x}$
 $= -x + (.1)\dot{x}^{3} - \beta\dot{x} + (.1)(3x^{2}\dot{x}) + d\dot{x}$
 $\dot{x} + x = (\alpha - \beta)\dot{x} + (.1)(\dot{x}^{3} + 3x^{2}\dot{x})$
3b. Harmonic balance: $x = A \cos \omega t$
 $-\omega^{2}A \cos \omega t + A \cos \omega t = (\alpha - \beta)(-A\omega \sin \omega t)$
 $+ (.1)(A^{3})(-\omega^{3}\sin^{3}\omega t)$
 $-3\omega^{3}\omega t \sin \omega t - \frac{1}{4}\sin^{3}\omega t$
 $\cos^{3}\omega t \sin \omega t = \frac{1}{4}\sin^{3}\omega t$
 $\cos^{3}\omega t \sin \omega t = \frac{1}{4}\sin^{3}\omega t$
 $\cos^{3}\omega t \sin \omega t = \frac{1}{4}\sin^{3}\omega t$
 $\cos^{3}\omega t \sin^{3}\omega t = \frac{1}{4}\sin^{3}(-\frac{1}{4})(-\frac{1$

3c. A Hopf bifurcation occurs for B=x] The LC exists for B7d. What is the stability of the origin for B7d? The DE becomes (Incarize hear the orgin); X+X = (d-13) × + nonlinion ferms for Bra this I is a damping term The origin is statle => the LC is unstable: Motions hear the LC more away from it and head towards the origin ... The Hopot beforeation is SUBCRITICAL (This agrees with simulation using pplane, Under "OPTIONS" Choose "SOLUTION DIRECTION" as forward. You will see the LC as repelling (i.e. unstable))