

①	KE (T)	PE (V)
Wheel I	$\frac{1}{2} I \dot{\theta}^2$	0
Torsional Spring	0	$\frac{1}{2} k_1 \theta^2$
Linear Spring	0	$\frac{1}{2} k_2 (\theta r - x)^2$
Mass M	$\frac{1}{2} M \dot{x}^2$	<del><math>M g x \sin 45</math></del> $\rightarrow 0$

The potential energy change of mass M due to gravity balances with the potential energy change due to static deflections at equilibrium, so  $V_M = 0$ .

Displacement of  $k_2 = \theta \cdot r - x = \Delta X_{k_2}$

y-displacement of M =  $x \sin 45 = \Delta y_M$

$$L = T - V$$

$$u_i = \{x, \theta\}$$

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{u}_i} \right) - \frac{\partial L}{\partial u_i} = 0$$

$$L = \frac{1}{2} I \dot{\theta}^2 + \frac{1}{2} M \dot{x}^2 - \frac{1}{2} k_1 \theta^2 - \frac{1}{2} k_2 (\theta r - x)^2$$

$$1) M \ddot{x} + k_2 (\theta r - x) (-1) = 0$$

$$2) I \ddot{\theta} + k_1 \theta + k_2 (\theta r - x) r = 0$$

$$1) M \ddot{x} - k_2 (\theta r - x) = 0$$

$$2) I \ddot{\theta} + k_1 \theta + k_2 r^2 \theta - k_2 r x = 0$$

$$\textcircled{2} \quad L = T - V$$

$$F = \begin{bmatrix} 0 \\ F \cos \Omega t \end{bmatrix}$$

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{x}_i} \right) - \frac{\partial L}{\partial x_i} = F_i$$

$$T = \frac{1}{2} m_1 \dot{x}_1^2 + \frac{1}{2} m_2 \dot{x}_2^2 = \frac{1}{2} \dot{x}_1^2 + \dot{x}_2^2$$

$$V = \frac{1}{2} k_1 x_1^2 + \frac{1}{2} k_2 (x_2 - x_1)^2 = \frac{1}{2} x_1^2 + (x_2^2 - 2x_2 x_1 + x_1^2)$$

$$= x_2^2 - 2x_2 x_1 + \frac{3}{2} x_1^2$$

$$L = T - V = \frac{1}{2} \dot{x}_1^2 + \dot{x}_2^2 - x_2^2 + 2x_1 x_2 - \frac{3}{2} x_1^2$$

$$1) \quad \ddot{x}_1 - 2x_2 + 3x_1 = 0$$

$$2) \quad 2\ddot{x}_2 + 2x_2 - 2x_1 = F \cos \Omega t$$

$$\underbrace{\begin{bmatrix} 1 & 0 \\ 0 & 2 \end{bmatrix}}_M \underbrace{\begin{bmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{bmatrix}} + \underbrace{\begin{bmatrix} 3 & -2 \\ -2 & 2 \end{bmatrix}}_K \underbrace{\begin{bmatrix} x_1 \\ x_2 \end{bmatrix}} = \underbrace{\begin{bmatrix} 0 \\ F \cos \Omega t \end{bmatrix}}_F$$

Take solution of form (at steady state):

$$x_1 = A F \cos \Omega t \quad x_2 = B F \cos \Omega t \quad \text{Want to minimize } |x_2|$$

$$1) \quad -A \Omega^2 - 2B + 3A = 0 \quad \rightarrow \quad A(3 - \Omega^2) = 2B$$

$$2) \quad -2B \Omega^2 - 2A + 2B = 1 \quad \rightarrow \quad 2B(1 - \Omega^2) - 2A = 1$$

$$2B(1 - \Omega^2) - \frac{4B}{3 - \Omega^2} = 1$$

$$B \left( 2(1 - \Omega^2) - \frac{4}{3 - \Omega^2} \right) = 1 \quad \rightarrow$$

2 cont'd

$$B = \frac{3 - \Omega^2}{(2)(1 - \Omega^2)(3 - \Omega^2) - 4}$$

Min  $|X_2|$  when  $B \rightarrow 0$

$B \rightarrow 0$  when  $(3 - \Omega^2) \rightarrow 0$   
 $\Omega^2 = 3$

$$B_{\Omega^2=3} = \frac{0}{-4} = 0$$

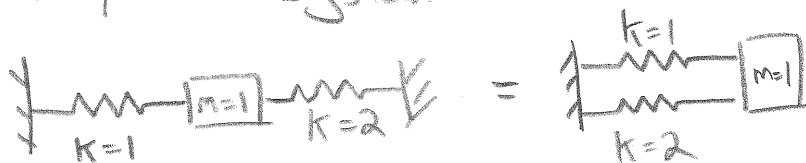
$\Omega = \sqrt{3}$  gives smallest  
steady state response  
of  $X_2$

Alternate Solution:

Assume smallest amplitude of  $X_2$  is when  $X_2$  is not moving

In this mode of vibration, would be driving the simplified system at its natural frequency, as motion of  $m_1$  will be  $180^\circ$  out of phase with the forcing.

Simplified system:



$$\omega_n = \Omega = \sqrt{\frac{k_{\text{eff}}}{m}} \quad k_{\text{eff}} = 2 + 1$$

$$\Omega = \sqrt{3}$$

③<sup>a</sup>  $\ddot{x} + 2n\dot{x} + x = 0$  Derivation of coefficients if you didn't know them!

General form (not knowing coefficients  $B$  or  $\omega_0$ )

$$x(t) = A e^{-Bt} \cos(\omega_0 t + \phi)$$

$$-A B e^{-Bt} \cos(\omega_0 t + \phi) - A \omega_0 e^{-Bt} \sin(\omega_0 t + \phi) = \dot{x}$$

$$+ A B^2 e^{-Bt} \cos(\omega_0 t + \phi) + A B \omega_0 e^{-Bt} \sin(\omega_0 t + \phi)$$

$$+ A B \omega_0 e^{-Bt} \sin(\omega_0 t + \phi)$$

$$- A \omega_0^2 e^{-Bt} \cos(\omega_0 t + \phi) = \ddot{x}$$

To save space,  $S = \sin(\omega_0 t + \phi) \cdot e^{-Bt}$

$C = \cos(\omega_0 t + \phi) \cdot e^{-Bt}$

Plug into Diff Eqn

$$A B^2 C + 2 A B \omega_0 S - A \omega_0^2 C + 2n(-A B C - A \omega_0 S) + A C = 0$$

$$C(A B^2 - A \omega_0^2 - 2n A B + A) = 0$$

$$C A (B^2 - \omega_0^2 - 2n B + 1) = 0 \Rightarrow B^2 - \omega_0^2 - 2n B + 1 = 0$$

$A \neq 0$ , uninteresting solution

$$S(2 A B \omega_0 - A \omega_0 \cdot 2n) = 0 \Rightarrow 2 B \omega_0 - 2n \omega_0 = 0 \Rightarrow B = n$$

$$n^2 - \omega_0^2 - 2n^2 + 1 = 0$$

$$\omega_0 = \sqrt{1 - n^2}$$

$$\Rightarrow x(t) = A e^{-nt} \cos(\sqrt{1 - n^2} t + \phi)$$

$$(3) \textcircled{a} \ddot{x} + 2n\dot{x} + x = 0$$

General form for damped harmonic motion:

$$\ddot{x} + 2\xi\omega_n\dot{x} + \omega_n^2 x = 0$$

$$x = Ae^{-\xi\omega_n t} \cos(\sqrt{1-\xi^2}\omega_n t + \phi)$$

$$\xi\omega_n = n \quad \omega_n^2 = 1 \Rightarrow \xi = n$$

$$x(t) = Ae^{-nt} \cos(\sqrt{1-n^2}t + \phi) \quad \leftarrow \text{Same as obtained if you didn't know general form coeffs}$$

$$x(t=0) = 0$$

$$Ae^{\phi} \cos \phi = 0$$

$$\cos \phi = 0 \quad \phi = \frac{\pi}{2}$$

$$x(t) = Ae^{-nt} \cos\left(\sqrt{1-n^2}t + \frac{\pi}{2}\right)$$

$$(b) \frac{x_1}{x_3} = 2 \quad \leftarrow \text{over two periods}$$

logarithmic decrement

$$\frac{x_1}{x_3} = \frac{e^{-n \cdot t_0}}{e^{-n \cdot (t_0 + 2T)}}$$

$$T = \frac{2\pi}{\omega_d} \quad \omega_d = \sqrt{1-n^2}$$

$$2 = e^{-n \cdot t_0 + n \cdot (t_0 + \frac{4\pi}{\sqrt{1-n^2}})}$$

$$\ln(2) = \frac{n \cdot 4\pi}{\sqrt{1-n^2}}$$

$$n^2 \cdot \frac{16\pi^2}{(\ln(2))^2} - 1 + n^2 = 0$$

$$n^2 \left( \frac{16\pi^2}{(\ln(2))^2} + 1 \right) = 1$$

$$n = \sqrt{\frac{1}{\frac{16\pi^2}{(\ln(2))^2} + 1}}$$



BC: longitudinal vibrations of a bar  $u_{tt} = c^2 u_{xx}$  where  $c^2 = \frac{EA}{\rho}$

1)  $u_x|_{x=l} = 0$  [free]

where  $\rho = \frac{\text{mass}}{\text{length}}$

2)  $u|_{x=0} = 0$  [clamped]

Assume a solution of the form

$u(x,t) = U_n(x) \cos(\omega_n t)$  plug into  $u_{tt} = c^2 u_{xx}$

$- \omega_n^2 U_n(x) \cos(\omega_n t) = c^2 U_n''(x) \cos(\omega_n t)$

$U_n''(x) + \frac{\omega_n^2}{c^2} U_n(x) = 0$

$U_n(x) = a \cos\left(\frac{\omega_n}{c} x\right) + b \sin\left(\frac{\omega_n}{c} x\right)$

To satisfy BC,  $U_n(x)$  must satisfy BC's as BC's are true for all  $t$

BC 1)  $-a \frac{\omega_n}{c} \sin\left(\frac{\omega_n}{c} x\right) + b \frac{\omega_n}{c} \cos\left(\frac{\omega_n}{c} x\right) = 0$  at  $x=l$

$\rightarrow -a \sin\left(\frac{\omega_n}{c} l\right) + b \cos\left(\frac{\omega_n}{c} l\right) = 0$

BC 2)  $a \cos(0) + b \sin(0) = 0$

$a = 0$

$\rightarrow b \cos\left(\frac{\omega_n}{c} l\right) = 0$   $b \neq 0$  (otherwise trivial soln)

$\frac{\omega_n}{c} l$  must be  $\frac{\pi}{2}, \frac{3\pi}{2}, \text{etc}$   $\frac{\omega_n}{c} l = \frac{(2n+1)\pi}{2}$

$n = 0, 1, 2, \dots$

1 continued

$$\omega_n = \frac{(2n+1)\pi}{2} \frac{c}{l}$$

$$\omega_n = \frac{(2n+1)\pi}{2} \frac{1}{l} \sqrt{\frac{EA}{\rho_{\text{length}}}} \quad n=0, 1, 2, \dots$$

②  $Q = \frac{EA}{\rho_{\text{vol}} A} \frac{\int_0^l V'(x)^2 dx}{\int_0^l V(x)^2 dx}$       Where  $\rho_{\text{vol}} = \frac{\text{mass}}{\text{Volume}}$   
since  $A$  is constant      (Note  $\rho_{\text{vol}} = \frac{\rho_{\text{length}}}{A}$ )

$$V(x) = a + bx + cx^2$$

BC's: longitudinal vibrations of a bar

$$V|_l = 0 \rightarrow b + 2cx = 0 \rightarrow b + 2cl = 0 \rightarrow b = -2cl$$

$$V|_0 = 0 \rightarrow a = 0$$

$$V(x) = c(-2lx + x^2)$$

$$V'(x) = c(-2l + 2x)$$

[Note: can take  $c=1$   
as it cancels out in  
 $Q$  anyway]

$$\text{num: } \int_0^l V'(x)^2 dx = c^2 \int_0^l (4x^2 - 8lx + 4l^2) dx$$

$$= c^2 \left( \frac{4}{3}x^3 - 4lx^2 + 4l^2x \right) \Big|_0^l$$

$$= c^2 \left( \frac{4}{3}l^3 - 4l^3 + 4l^3 \right) = c^2 \left( \frac{4}{3}l^3 \right)$$

2 continued

$$\begin{aligned} \text{den: } \int_0^l V(x)^2 dx &= c^2 \int_0^l x^4 - 4lx^3 + 4l^2x^2 dx \\ &= c^2 \left( \frac{x^5}{5} - lx^4 + \frac{4}{3}l^2x^3 \right) \Big|_0^l \\ &= c^2 l^5 \left( \frac{1}{5} - 1 + \frac{4}{3} \right) \\ & \quad \left( \frac{1}{5} + \frac{1}{3} \right) \\ &= c^2 l^5 \left( \frac{8}{15} \right) \end{aligned}$$

$$Q = \frac{E}{\rho_{\text{vol}}} \frac{\cancel{4} \frac{4}{3} l^3}{\cancel{4} \frac{8}{15} l^5} = \frac{5}{2} \frac{E}{l^2 \rho_{\text{vol}}}$$

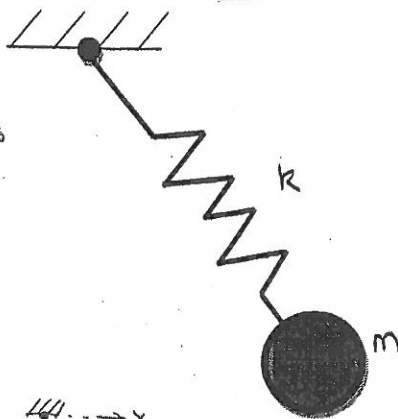
$$\frac{4}{3} \cdot \frac{15}{8} = \frac{5}{2}$$

$$W_1 \leq \sqrt{Q}$$

$$W_1 \leq \sqrt{\frac{5}{2} \frac{E}{\rho_{\text{vol}} l^2}}$$



★ "derive the model"  
find the equations of motion, clearly label/explain all variables & coordinates



$l_0 =$  unstretched length

• Cartesian Coordinates

$$T = \frac{1}{2} m \dot{x}^2 + \frac{1}{2} m \dot{y}^2$$

$$U = -mgy + \frac{1}{2} k (\sqrt{x^2 + y^2} - l_0)^2$$

$$L = T - U = \frac{1}{2} m \dot{x}^2 + \frac{1}{2} m \dot{y}^2 + mgy - \frac{1}{2} k (\sqrt{x^2 + y^2} - l_0)^2$$

$$\frac{d}{dt} \frac{\delta L}{\delta \dot{x}} - \frac{\delta L}{\delta x} = 0 \quad \& \quad \frac{d}{dt} \frac{\delta L}{\delta \dot{y}} - \frac{\delta L}{\delta y} = 0$$

$$\frac{\delta L}{\delta \dot{x}} = m \dot{x}$$

$$\frac{\delta L}{\delta \dot{y}} = m \dot{y}$$

$$\frac{d}{dt} \frac{\delta L}{\delta \dot{x}} = m \ddot{x}$$

$$\frac{d}{dt} \frac{\delta L}{\delta \dot{y}} = m \ddot{y}$$

$$\frac{\delta L}{\delta x} = -kx + \frac{k l_0 x}{\sqrt{x^2 + y^2}}$$

$$\frac{\delta L}{\delta y} = mg - ky + \frac{k l_0 y}{\sqrt{x^2 + y^2}}$$

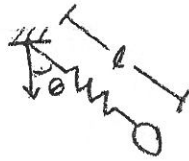
so equations of motion in Cartesian coordinates are

$$m \ddot{x} + kx - \frac{k l_0 x}{\sqrt{x^2 + y^2}} = 0$$

$$\& \quad m \ddot{y} - mg + ky - \frac{k l_0 y}{\sqrt{x^2 + y^2}}$$

(1)

~~Cartesian Coordinates~~  
 • Polar Coordinates



$$x = l \sin \theta \quad \dot{x} = \dot{l} \sin \theta + l \dot{\theta} \cos \theta$$

$$y = l \cos \theta \quad \dot{y} = \dot{l} \cos \theta - l \dot{\theta} \sin \theta$$

convert  $T$  &  $U$  from cartesian to polar coordinates:

$$T = \frac{1}{2} m \dot{x}^2 + \frac{1}{2} m \dot{y}^2$$

$$= \frac{1}{2} m [(\dot{l} \sin \theta + l \dot{\theta} \cos \theta)^2 + (\dot{l} \cos \theta - l \dot{\theta} \sin \theta)^2]$$

$$= \frac{1}{2} m [\dot{l}^2 \sin^2 \theta + 2 \dot{l} l \dot{\theta} \cos \theta \sin \theta + l^2 \dot{\theta}^2 \cos^2 \theta + \dot{l}^2 \cos^2 \theta - 2 \dot{l} l \dot{\theta} \cos \theta \sin \theta + l^2 \dot{\theta}^2 \sin^2 \theta]$$

$$T = \frac{1}{2} m \dot{l}^2 + \frac{1}{2} m l^2 \dot{\theta}^2$$

$$U = -mgy + \frac{1}{2} k (\sqrt{x^2 + y^2} - l_0)^2$$

$$= -mgl \cos \theta + \frac{1}{2} k (\sqrt{l^2 \sin^2 \theta + l^2 \cos^2 \theta} - l_0)^2$$

$$U = -mgl \cos \theta + \frac{1}{2} k (l - l_0)^2$$

$$L = \frac{1}{2} m \dot{l}^2 + \frac{1}{2} m l^2 \dot{\theta}^2 + mgl \cos \theta - \frac{1}{2} k (l - l_0)^2$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}} - \frac{\partial L}{\partial \theta} = 0 \quad \& \quad \frac{d}{dt} \frac{\partial L}{\partial \dot{l}} - \frac{\partial L}{\partial l} = 0$$

$$\frac{\partial L}{\partial \dot{\theta}} = m l^2 \dot{\theta} \quad \frac{\partial L}{\partial \dot{l}} = m \dot{l}$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}} = m l^2 \ddot{\theta} + 2 m l \dot{l} \dot{\theta} \quad \frac{d}{dt} \frac{\partial L}{\partial \dot{l}} = m \ddot{l}$$

$$\frac{\partial L}{\partial \theta} = -mgl \sin \theta \quad \frac{\partial L}{\partial l} = m l \dot{\theta}^2 + mg \cos \theta - k(l - l_0)$$

so equations of motion in polar coordinates are

$$l \ddot{\theta} + 2 \dot{l} \dot{\theta} + g \sin \theta = 0 \quad \& \quad m \ddot{l} - m l \dot{\theta}^2 + k(l - l_0) - mg \cos \theta = 0$$