

Solution - Basant Sharma Mar 31<sup>st</sup> 02

① 7.1

The trajectory of given particle is

$$\vec{r}(t) = R \cos(\dot{\theta}t) \hat{i} + R \sin(\dot{\theta}t) \hat{j}$$

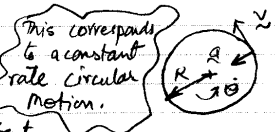
where  $\dot{\theta} = \text{constant}$ .

so  $\vec{v}(t) = \dot{\vec{r}}(t) = R\dot{\theta}(-\sin(\dot{\theta}t)\hat{i} + \cos(\dot{\theta}t)\hat{j})$

a)  $\vec{v} \neq \vec{a}$  as  $\vec{v}$  depends on time.

(take  $t=0$ , for example)

b)  $\vec{v} \neq \text{constant}$  (see a)



c)  $|\vec{v}| = R\dot{\theta} = \text{constant}$

Since R and  $\dot{\theta}$  are constant.

d)  $\vec{a} = \dot{\vec{v}} = -R\dot{\theta}^2(\cos(\dot{\theta}t)\hat{i} + \sin(\dot{\theta}t)\hat{j})$

$\neq 0$  (take  $t=0$  for example)

e)  $\vec{a} \neq \text{constant}$  as it is clearly a function of time t.

f)  $|\vec{a}| = R\dot{\theta}^2 = \text{constant}$  (see c)

g)  $\vec{v} \cdot \vec{a} = -R^2\dot{\theta}^3(-\sin(\dot{\theta}t)\cos(\dot{\theta}t) + \cos(\dot{\theta}t)\sin(\dot{\theta}t)) = 0$

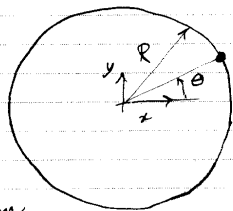
$\therefore \vec{v} \perp \vec{a}$ .

② 7.8

Bead on a circular path

The rate of change of angular speed is proportional to the angular position as  $\dot{\omega} = c\omega^{3/2}$

So  $\ddot{\theta} = c\theta^{3/2}$  (1)



Given the bead starts from rest at  $\theta=0$ ,

So  $\theta(0)=0, \dot{\theta}(0)=0$ . One could guess that  $\theta=0$  is the solution. Let's do it more rigorously. Multiplying both sides by  $\dot{\theta}$  (say  $\dot{\theta} \neq 0$ ),  $\dot{\theta}\ddot{\theta} = c\theta^{3/2}\dot{\theta}$

or  $\frac{1}{2} \frac{d\dot{\theta}^2}{dt} = \frac{c}{5/2} \frac{d\theta^{5/2}}{dt}$  (2) (Use "inverse" chain rule)

Integrating (2) w.r.t. t we get

$$\frac{\dot{\theta}^2}{2} = \frac{2c}{5} \theta^{5/2} + C_1$$

at  $t=0, \theta=\dot{\theta}=0 \Rightarrow C_1=0$

So  $\dot{\theta}^2 = \frac{4c}{5} \theta^{5/2}$

Consider only forward motion,  $\dot{\theta} = \sqrt{\frac{4c}{5}} \theta^{5/4}$

Using separation of variables

$$\theta^{5/4} = -1/4 \sqrt{\frac{4c}{5} t + C_2}$$

Since  $t=0, \theta=0$  so  $C_2 = \infty!$

So  $\dot{\theta} \neq 0$  assumption was not good. So

Only solution is  $\theta=0$ .

Solution after Correction:

Here  $\ddot{\theta} = c\theta^{3/2}$  (3)

Assume  $\dot{\theta} \neq 0$  then  $\frac{\dot{\theta}^2}{2} = \frac{c}{3/2} \theta^{3/2} + C_1$

(Same reasoning as above)

Since at  $t=0, \theta=\dot{\theta}=0, C_1=0$

So  $\dot{\theta} = \sqrt{\frac{4c}{3}} \theta^{3/4}$

By separation of variables and using  $\theta=0$  at  $t=0$

$$\theta^{1/4} = \frac{1}{4} \sqrt{\frac{4c}{3}} t + C_2 \Rightarrow \theta^{1/4} = \frac{1}{4} \sqrt{\frac{4c}{3}} t + C_2 \Rightarrow \theta = \left(\frac{1}{4} \sqrt{\frac{4c}{3}} t + C_2\right)^4$$

So  $\theta = \left(\frac{1}{4} \sqrt{\frac{4c}{3}} t\right)^4 = \frac{c^2}{144} t^4$

Other solution is clearly  $\dot{\theta} \equiv 0$  i.e.  $\theta \equiv 0$

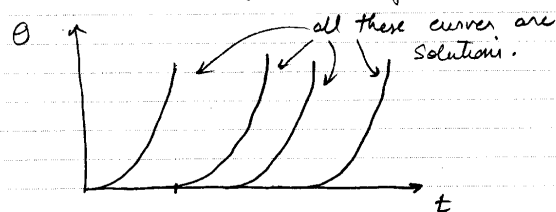
Since  $\theta=0$  at  $t=0$ .

So we have clearly two distinct solutions to (3) with given initial condition.

In fact there are infinite number of solutions and each such solution is given by

$$\theta = \begin{cases} 0 & 0 \leq t \leq t_1 \\ \frac{c^2}{144} t^4 & t_1 \leq t \end{cases}$$

where  $t_1$  is any number of choice!



Note: That an ODE has more than one solution with given ICs (like this) is exceptional (you won't see it too often).

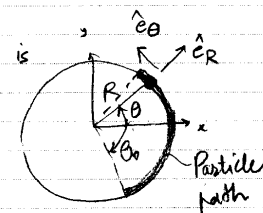
③ 7.14 We have to find the conditions for which  $|\vec{a}|$  is max at  $\theta=0$ .

Given  $\theta = \theta_0 \cos \lambda t$

So  $\dot{\theta} = -\theta_0 \lambda \sin \lambda t$

$\ddot{\theta} = -\theta_0 \lambda^2 \cos \lambda t$

And  $\vec{a} = -R\ddot{\theta} \hat{e}_R + R\dot{\theta} \dot{\theta} \hat{e}_\theta$



where  $\dot{\theta}^2 = \theta_0^2 \lambda^2 \sin^2 \lambda t = \theta_0^2 \lambda^2 (1 - \cos^2 \lambda t) = \lambda^2 (\theta_0^2 - \theta^2); \ddot{\theta} = -\lambda^2 \theta$

So  $|\vec{a}|^2 = R^2 \dot{\theta}^4 + R^2 \ddot{\theta}^2$

$= R^2 (\lambda^4 (\theta_0^2 - \theta^2)^2 + \lambda^4 \theta^2)$

$= R^2 \lambda^4 ((\theta_0^2 - \theta^2)^2 + \theta^2)$  (all f(\theta))

explicit Note: time dependence is eliminated!

Maximum acceleration occurs at the maximum value of  $f(\theta) = \theta_0^4 + \theta^4 + \theta^2(1-2\theta^2)$  as  $\lambda, R$  are given constants.

at  $\theta=0, f(\theta) = \theta_0^4$

So we need  $f(\theta) \leq \theta_0^4$  for any  $\theta$

i.e.  $\theta^4 + \theta^2(1-2\theta^2) \leq \theta_0^4 \Rightarrow \theta^2 + 1 - 2\theta^2 \leq \theta_0^2$  (as  $\theta^2 \geq 0$ )

We know that  $\theta = \theta_0 \cos \lambda t$

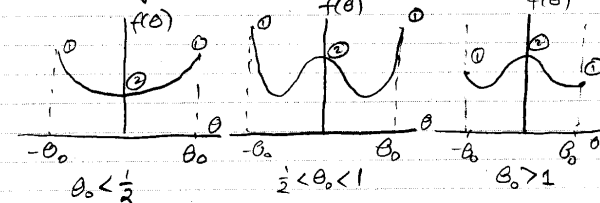
so  $|\theta| \leq |\theta_0|$  or  $\theta^2 \leq \theta_0^2$

Therefore, for maximum acceleration we need  $1 - \theta_0^2 \leq 0$  or  $\theta_0 \geq 1$

Then the maximum acceleration becomes

$|\vec{a}| = R\lambda^2 \theta_0^2$

Graphically:  $f(\theta) = (\theta_0^2 - \theta^2)^2 + \theta^2 = \theta^4 + \theta^2(1-2\theta^2) + \theta_0^4$



①  $f(\theta)|_{\theta=\theta_0} = \theta_0^2$ ; ②  $f(\theta=0) = \theta_0^4$

$\therefore$  Max value of  $|\vec{a}|$  will be at  $\theta=0$  only when  $\theta_0 \geq 1 \Rightarrow |\vec{a}|_{\text{max}} |\vec{a}|_{\theta=0} = R\lambda^2 \theta_0^2$ .

④ 7.41 Simple Pendulum,  $M=1\text{ kg}$ ,  $L=1\text{ m}$   
(no air friction)  $g=10\text{ m/s}^2$

a) LMB:  $\Sigma \vec{F} = m\vec{a}$

$$\{-T\hat{e}_R + mg\hat{i} = -mL\dot{\theta}^2\hat{e}_R + mL\ddot{\theta}\hat{e}_\theta\}$$

$$\{ \} \cdot \hat{e}_\theta \Rightarrow mg(\hat{i} \cdot \hat{e}_\theta) = mL\ddot{\theta}$$

where  $\hat{i} \cdot \hat{e}_\theta = \hat{i} \cdot (-\sin\theta\hat{i} + \cos\theta\hat{j})$

$$= -\sin\theta$$

$$\therefore -mg\sin\theta = mL\ddot{\theta}$$

$$\Rightarrow \ddot{\theta} + \frac{g}{L}\sin\theta = 0 \quad (1)$$

b) LMB.  $\hat{e}_R \Rightarrow -T + mg(\hat{i} \cdot \hat{e}_R) = -mL\dot{\theta}^2$

where  $\hat{i} \cdot \hat{e}_R = \hat{i} \cdot (\cos\theta\hat{i} + \sin\theta\hat{j}) = \cos\theta$

$$\Rightarrow T = mg\cos\theta + mL\dot{\theta}^2 \quad (2)$$

c) The force at the hinge support is simple  $\vec{F}_0 = -T\hat{e}_R$  since the string is massless.

$$\vec{F}_0 = -T\hat{e}_R = -T(\cos\theta\hat{i} + \sin\theta\hat{j})$$

$$\text{so } F_{0x}\hat{i} + F_{0y}\hat{j} = -(mg\cos\theta + mL\dot{\theta}^2)(\cos\theta\hat{i} + \sin\theta\hat{j})$$

$$\text{or } F_{0x} = -(mg\cos\theta + mL\dot{\theta}^2)\cos\theta; F_{0y} = -(mg\cos\theta + mL\dot{\theta}^2)\sin\theta$$

d) Let  $\omega = \dot{\theta} \Rightarrow \dot{\omega} = \ddot{\theta} = -\frac{g}{L}\sin\theta$  by (1)

$$\Rightarrow \begin{cases} \dot{\theta} = \omega \\ \dot{\omega} = -\frac{g}{L}\sin\theta \end{cases}$$

e) See next page.

f) Max. tension = 29.999 N

g) Time period = 2.3452 s ← by fine tuning tspan, interpolation, or event detection.

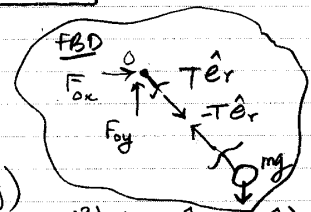
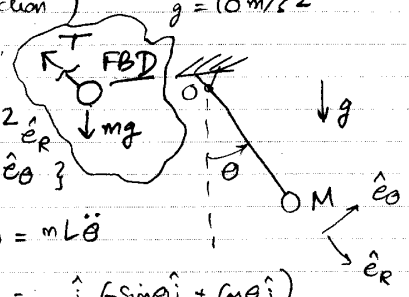
h) Note:

① Maximum Tension is independent of string length.

Since Energy (total) is conserved, initial

$$\frac{1}{2} m(L\dot{\theta})^2 + mg(L - L\cos\theta) = mgL$$

$$\text{ie. } mL\dot{\theta}^2 + 2mg(1 - \cos\theta) = 2mg$$



$$\text{So } mL\dot{\theta}^2 = 2mg\cos\theta$$

$$\text{So by (2) } T = 3mg\cos\theta$$

$$\Rightarrow T_{\text{max}} = 3mg$$

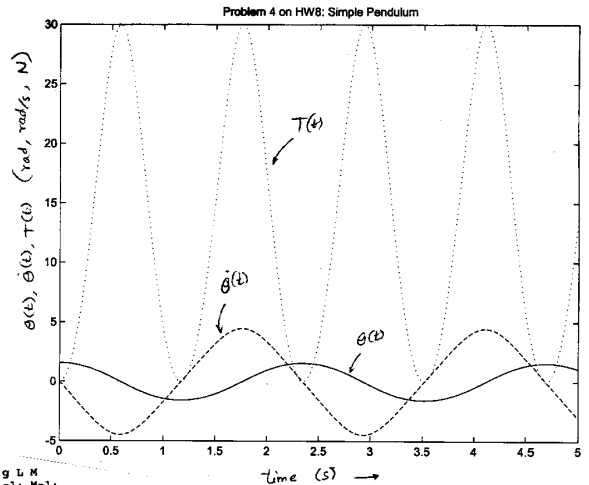
If you compute numerically  $T_{\text{max}}$  then you will always obtain value close to 30

L	$T_{\text{max}}$ (N)
1	29.9993
5	29.9935
10	29.9977

② Maximum tension can be made closer to 30N by decreasing the 'RelTol' and 'AbsTol'. For example:

AbsTol	RelTol	$T_{\text{max}}$
1e-6	1e-6	29.9993
1e-4	1e-4	29.9820
1e-4	1e-8	29.9947
1e-8	1e-8	30.0000

③ Use  $\sin\theta \approx \theta$  and get Time Period = 1.9869 sec (see part a)



```
global g L M
g=10; L=1; M=1;
tspan=[0 5];
theta0=pi/2; thetadot0=0;
z0=[theta0 thetadot0];
options=odeset('AbsTol',1e-6,'RelTol',1e-6);

function zdot=simplependulum(t,z)
global g L M
theta=z(1);
theadot=z(2);
omega=z(2);

theadot-omega;
omegadot=(-g/L)*sin(theta);
zdot=[theadot omegadot]';

[t,z]=ode45('simplependulum',tspan,z0,options);

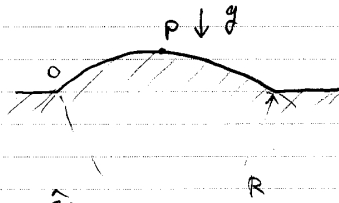
theta=z(:,1);
theadot=z(:,2);
tension=M*g*cos(theta)+M*L*theadot.^2;
max(tension)

plot(t,theta)
hold on
plot(t,theadot,'--')
plot(t,tension,':')

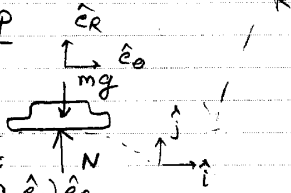
title('Problem 4 on HW8: Simple Pendulum');
```

⑤ 7.47

Car moves with speed  $v$  and climbs the hill at  $O$



FBD of car at P



By the Second law in the radial direction at P:

$$(N - mg) \hat{e}_R = m(\ddot{r} - r\dot{\theta}^2) \hat{e}_R$$

$$= m\left(-\frac{v^2}{R}\right) \hat{e}_R \quad (\because \dot{v} = 0)$$

If the car just barely leaves the ground at P then  $N = 0$

$$\Rightarrow mg = \frac{mv^2}{R} \quad \text{or} \quad \boxed{v = \sqrt{gR}}$$

⑥ 7.59

Circular disc  $R = 100 \text{ mm}$  rotating with center  $O$ .

$V_A = |V_A| = 0.8 \text{ m/s}$  at given instant.

$a_B$  makes angle  $\theta$  with radial line as shown.

Angular acceleration  $\alpha = ?$

Given  $\tan \theta = 0.6$ .

By kinematics

$$V = \omega \times r$$

$$a = \alpha \times r - \omega^2 r$$

where  $r = r \hat{e}_r$ ,  $\omega = \omega \hat{k}$ ,  $\alpha = \alpha \hat{k}$  with  $\omega$  and  $\alpha$  instantaneous angular velocity and acceleration respectively. Using  $\hat{k} \times \hat{e}_r = \hat{e}_\theta$  we get,

for point A:  $V_A = \omega R \Rightarrow \omega = \frac{V_A}{R}$  — (1)

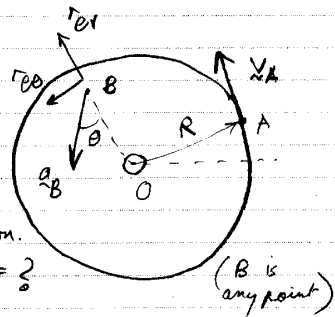
for point B:  $a_B = \alpha r \hat{e}_\theta - \omega^2 r \hat{e}_r$  — (2)

$$(2) \Rightarrow \tan \theta = \frac{a_B \cdot \hat{e}_\theta}{-(a_B \cdot \hat{e}_r)} = \frac{\alpha r}{\omega^2 r} = \frac{\alpha}{\omega^2}$$

by (1),  $\tan \theta = \frac{\alpha}{(V_A)^2 / R^2}$  0.6

Using the data given

$$\alpha = \frac{V_A^2 \tan \theta}{R^2} = 38.4 \frac{\text{rad}}{\text{s}^2}$$



⑦

Matlab Code - rotatepic.m

```
before=[1 2 1.5 1 1 2 2 1; 1 1 2 1 2 3 3 1];
line(before(1,:),before(2,:)); % input picture
axis equal;
hold on;
```

```
theta=2*pi/3;
Q=[cos(theta) -sin(theta); sin(theta) cos(theta)];
after=Q*before;
line(after(1,:),after(2,:)); % rotated picture
title('Rotation of a drawing by 120 deg');
```

